



Advanced Computation:
Computational Electromagnetics

Implementation of Finite-Difference Frequency-Domain

1

Outline

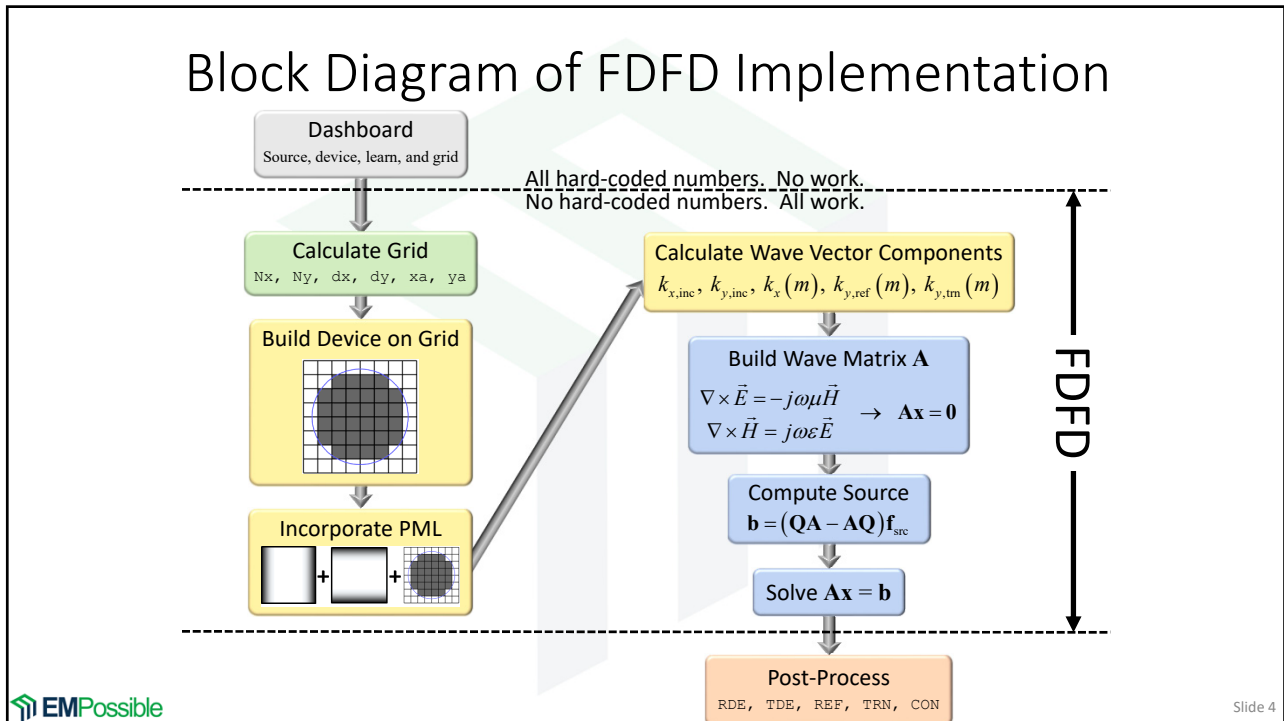
- Basic flow of FDFD
- 2× grid technique
- Calculating grid parameters
- Constructing your device on the grid
- Walkthrough of the FDFD algorithm
- Examples for benchmarking

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Basic Flow of FDFD

Slide 3

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Slide 4

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Detailed FDFD Algorithm

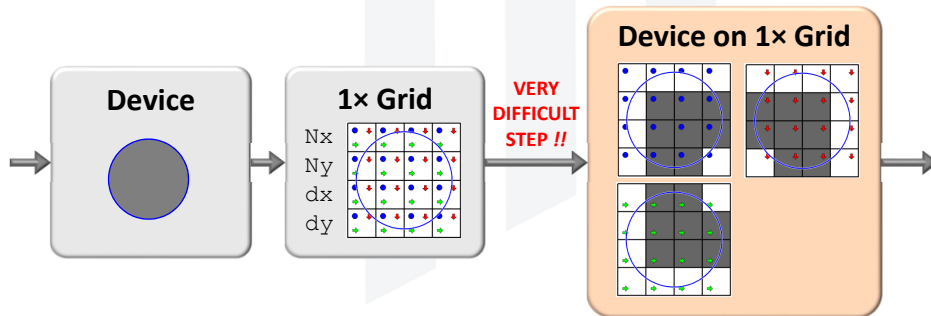
1. Construct FDFD Problem
 - a. Define your problem
 - b. Calculate the grid parameters
 - c. Assign materials to the grid to build ER2 and UR2 arrays
2. Handle PML and Materials
 - a. Compute s_x and s_y
 - b. Incorporate into ϵ_r and μ_r
 - c. Overlay onto 1x grids
3. Compute Wave Vector Components
 - a. Identify materials in reflected and transmitted regions: ϵ_{ref} , ϵ_{trn} , μ_{ref} , μ_{trn} , n_{ref} , n_{trn}
 - b. Compute incident wave vector: \mathbf{k}_{inc}
 - c. Compute transverse wave vector expansion: $k_{x,m}$
 - d. Compute $k_{y,ref}$ and $k_{y,trn}$
4. Construct **A**
 - a. Construct diagonal materials matrices
 - b. Compute derivative matrices
 - c. Compute **A**
5. Compute Source Vector, **b**
 - a. compute source field \mathbf{f}_{src}
 - b. compute **Q**
 - c. compute source vector **b**
6. Solve Matrix Problem: $\mathbf{e} = \mathbf{A}^{-1}\mathbf{b}$
7. Post Process Data
 - a. Extract E_{ref} and E_{trn}
 - b. Remove phase tilt
 - c. FFT the fields
 - d. Compute diffraction efficiencies
 - e. Compute reflectance & transmittance
 - f. Compute conservation of power

2x Grid Technique

What is the 2x Grid Technique? (1 of 2)

This is the traditional approach for building devices on a Yee grid.

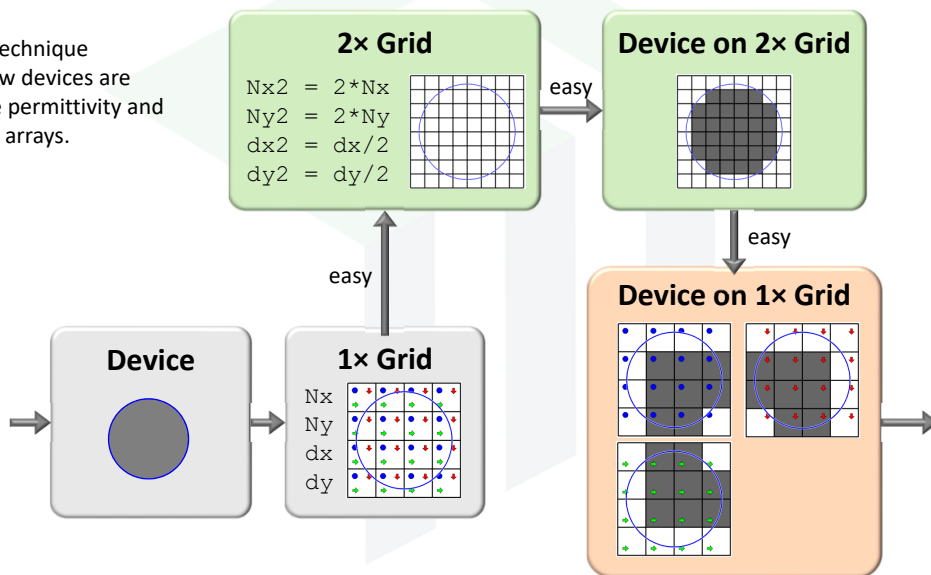
It is very tedious and cumbersome to determine which field components reside in which material.



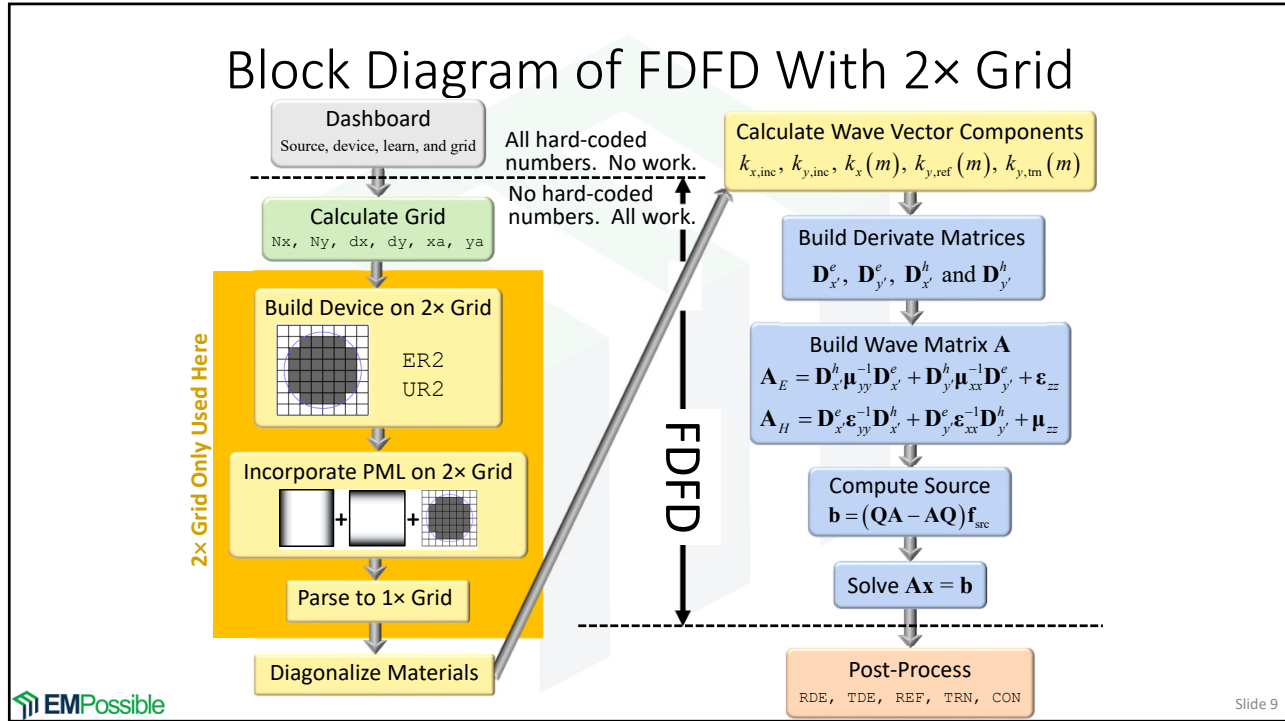
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What is the 2x Grid Technique? (2 of 2)

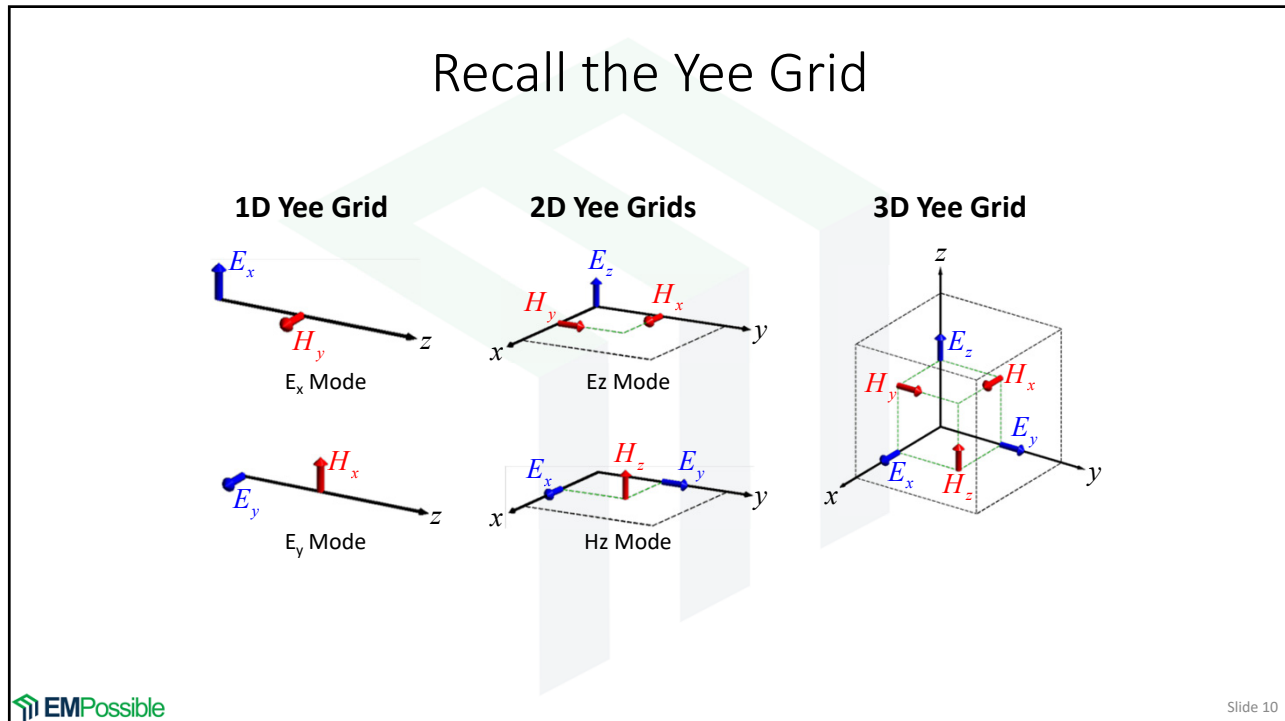
The 2x grid technique simplifies how devices are built into the permittivity and permeability arrays.



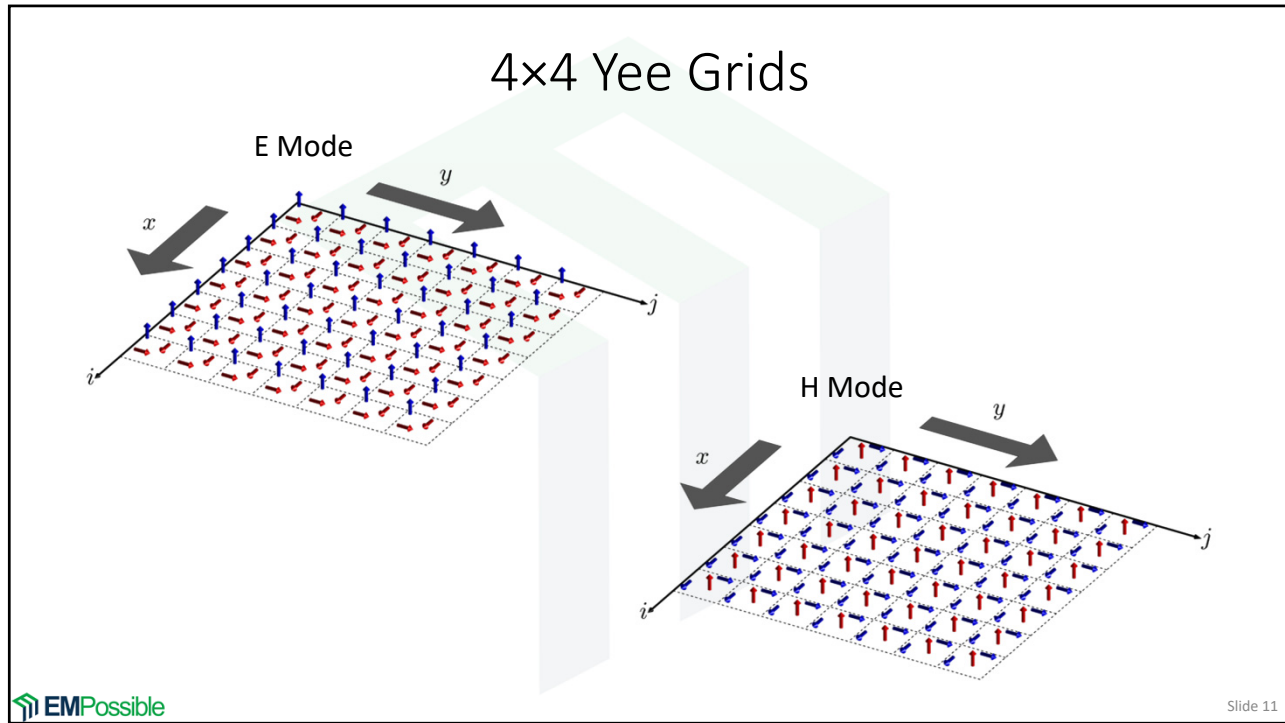
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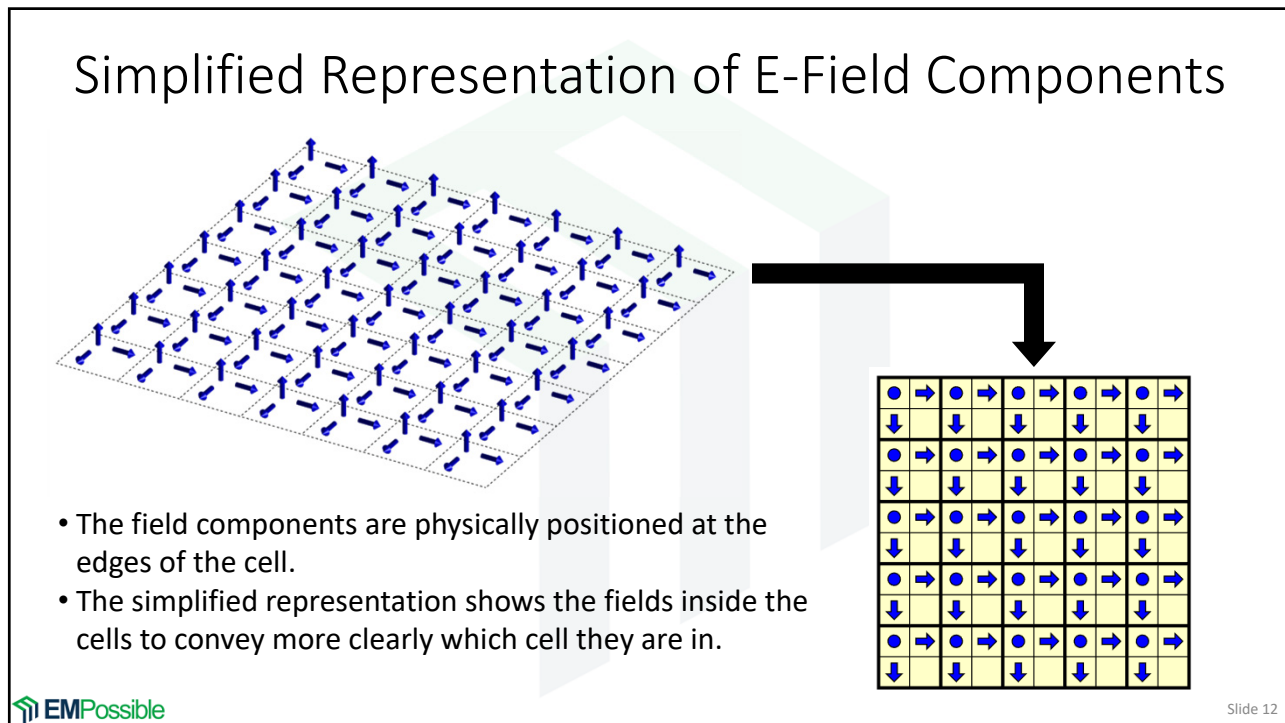
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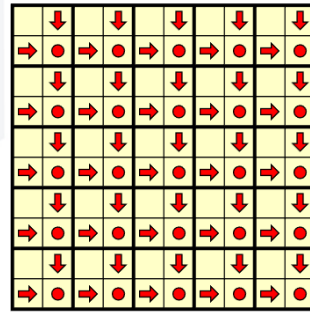
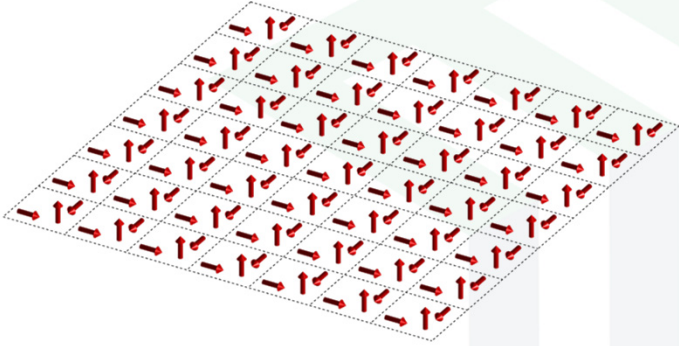


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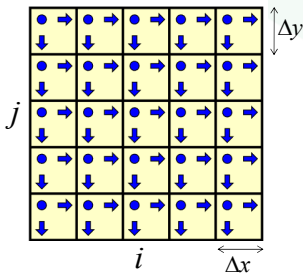
Simplified Representation of H-Field Components



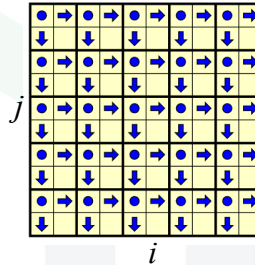
- The field components are physically positioned at the edges of the cell.
- The simplified representation shows the fields inside the cells to convey more clearly which cell they are in.

The 2x Grid

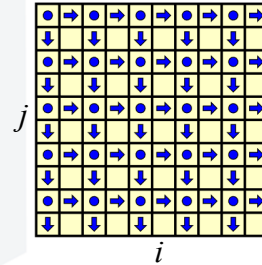
The Conventional 1x Grid



Due to the staggered nature of the Yee grid, we are effectively getting twice the resolution.



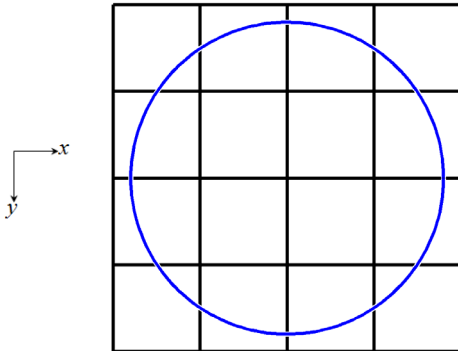
It now makes sense to talk about a grid that is at twice the resolution, the "2x grid."



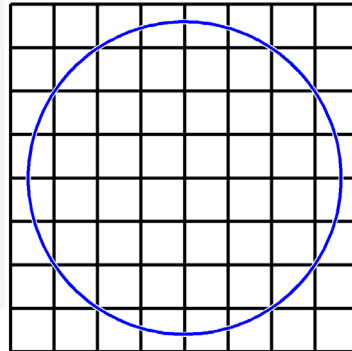
The 2x grid concept is useful because we can create devices (or PMLs) on the 2x grid without having to think about where the different field components are located. In a second step, we can easily pull off the values from the 2x grid where they exist for a particular field component.

$2\times \rightarrow 1\times$ (1 of 4): *Define Grids*

Suppose a circle is to be built onto the Yee grid. First, define the standard $1\times$ grid.



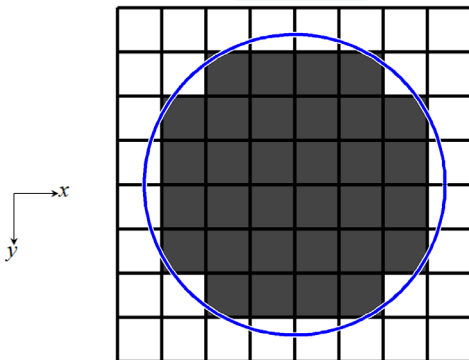
Second, define a corresponding $2\times$ grid.



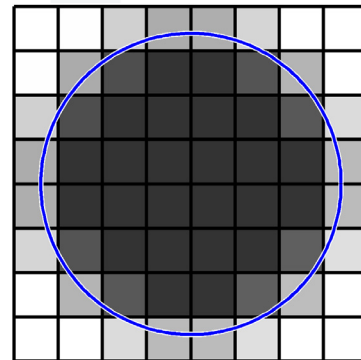
Note: the $2\times$ grid represents the same physical space, but with twice the number of points.

$2\times \rightarrow 1\times$ (2 of 4): *Build Device*

Third, we construct a cylinder on the $2\times$ grid without having to consider anything about the Yee grid.



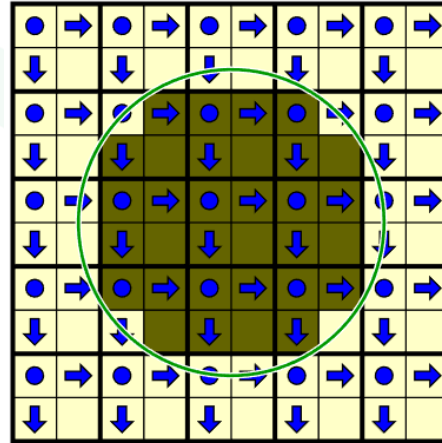
Fourth, if desired we could perform dielectric averaging on the $2\times$ grid at this point.



Note: This is discussed more in Lecture 14 FDFD Extras.

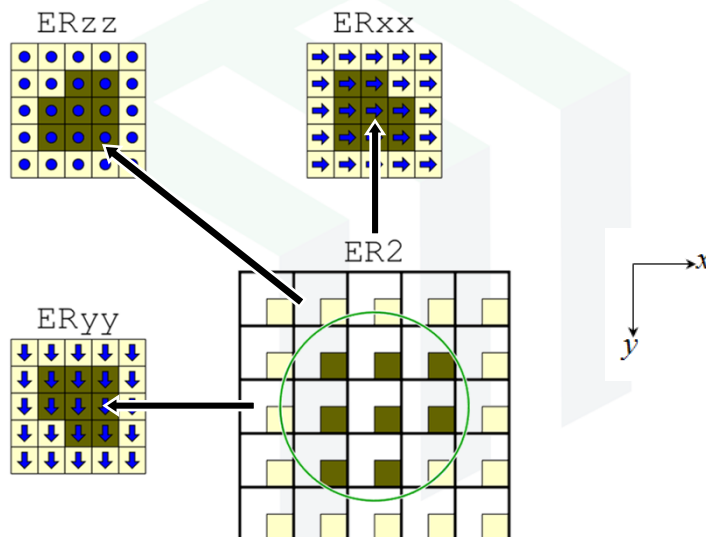
$2\times \rightarrow 1\times$ (3 of 4): Recall Field Staggering

Recall the relation between the $2\times$ and $1\times$ grids as well as the location of the field components.



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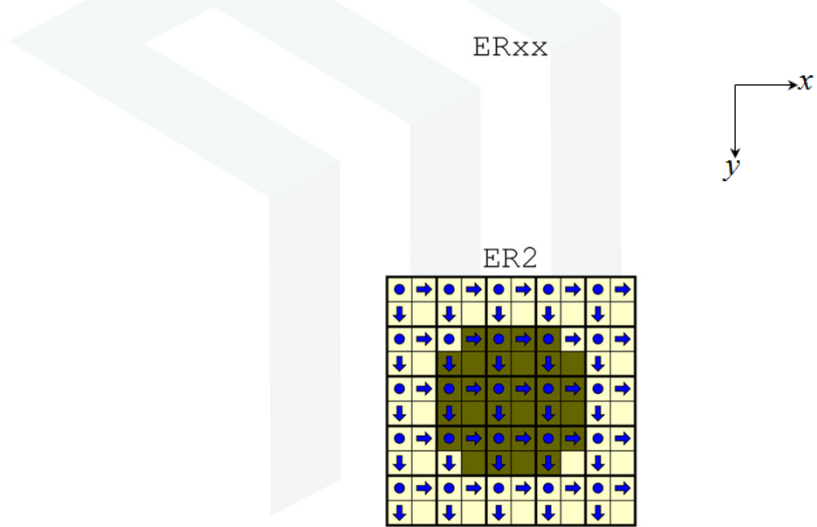
$2\times \rightarrow 1\times$ (4 of 4): Parse to $1\times$ Grid



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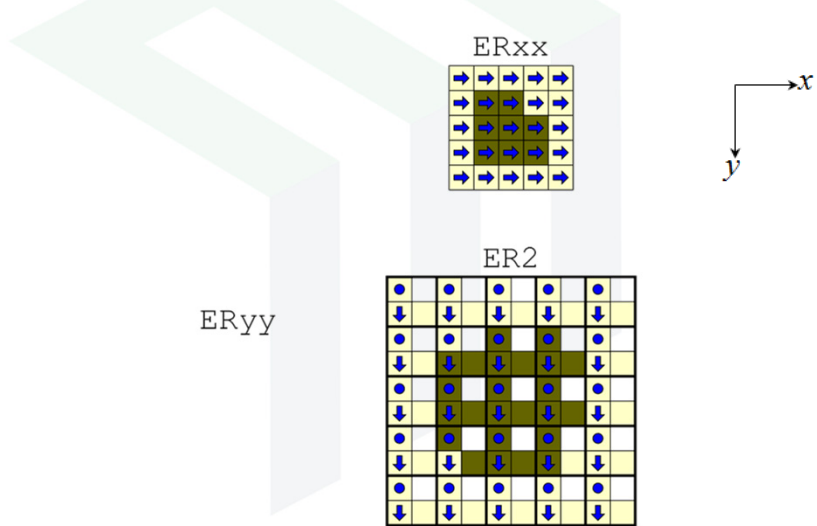
Extract ER_{xx} from ER₂

$$ER_{xx} = ER_2(2:2:Nx2, 1:2:Ny2)$$



Extract ER_{yy} from ER₂

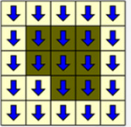
$$ER_{yy} = ER_2(1:2:Nx2, 2:2:Ny2)$$



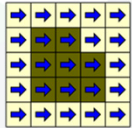
Extract ERzz from ER2

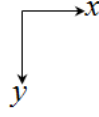
$$ER_{zz} = ER2(1:2:Nx2, 1:2:Ny2)$$

ERzz

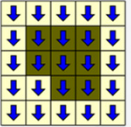


ERxx

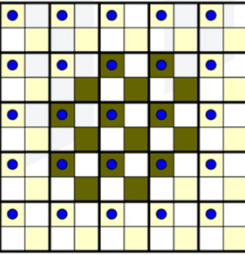





ERyy



ER2



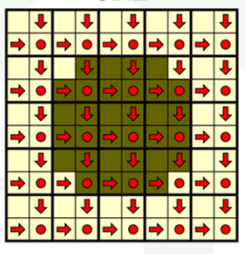

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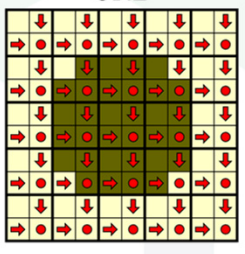
Extract URxx from UR2

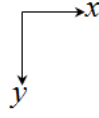
$$UR_{xx} = UR2(1:2:Nx2, 2:2:Ny2)$$


UR2



URxx



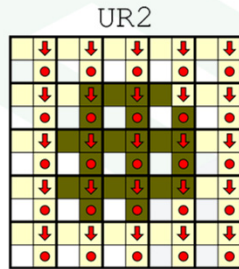



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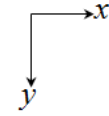
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Extract UR_{yy} from UR₂

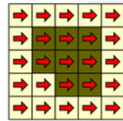
$$UR_{yy} = UR_2(2:2:Nx2, 1:2:Ny2)$$



UR_{yy}

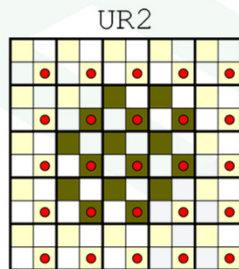


UR_{xx}

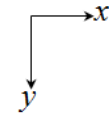
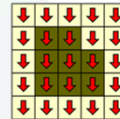


Extract UR_{zz} from UR₂

$$UR_{zz} = UR_2(2:2:Nx2, 2:2:Ny2)$$

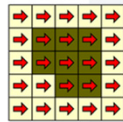


UR_{yy}

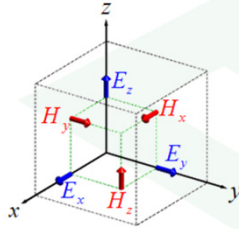


UR_{zz}

UR_{xx}

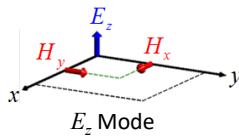


MATLAB Code for Parsing Onto 1x Grid

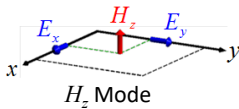


```
ERxx = ER2(2:2:Nx2, 1:2:Ny2, 1:2:Nz2);
ERyy = ER2(1:2:Nx2, 2:2:Ny2, 1:2:Nz2);
ERzz = ER2(1:2:Nx2, 1:2:Ny2, 2:2:Nz2);

URxx = UR2(1:2:Nx2, 2:2:Ny2, 2:2:Nz2);
URyy = UR2(2:2:Nx2, 1:2:Ny2, 2:2:Nz2);
URzz = UR2(2:2:Nx2, 2:2:Ny2, 1:2:Nz2);
```

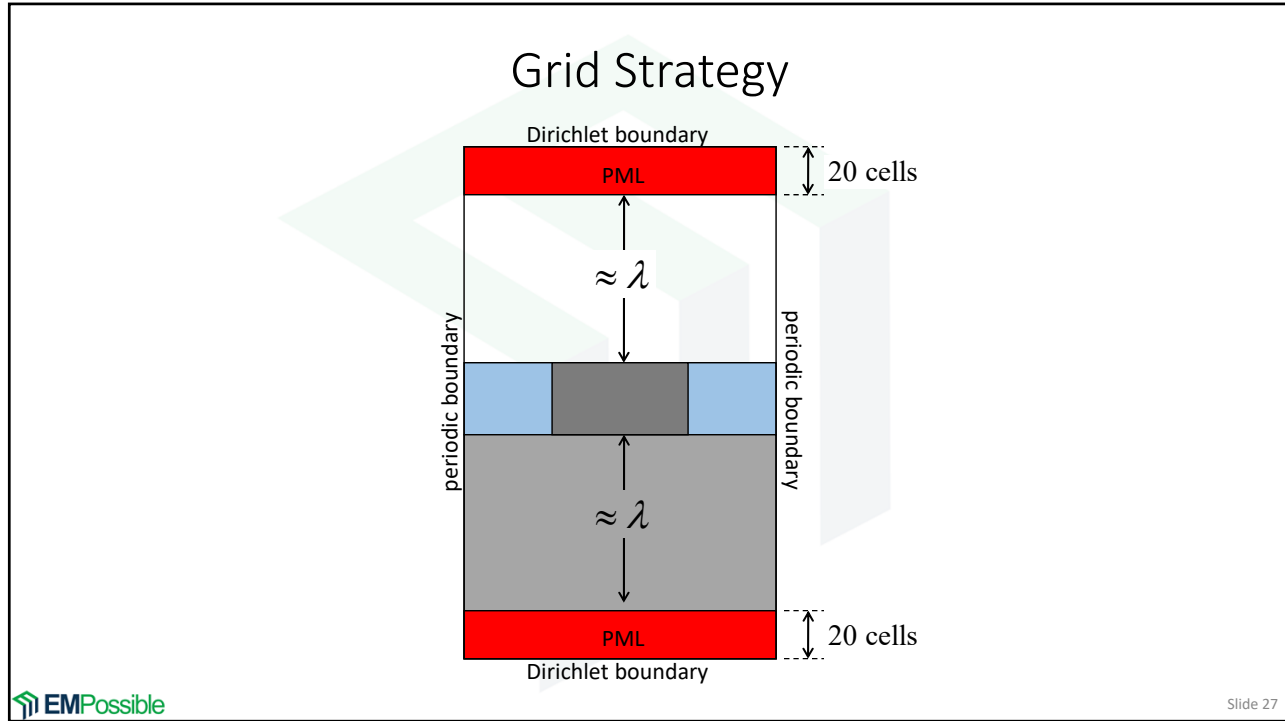


```
ERzz = ER2(1:2:Nx2, 1:2:Ny2);
URxx = UR2(1:2:Nx2, 2:2:Ny2);
URyy = UR2(2:2:Nx2, 1:2:Ny2);
```



```
ERxx = ER2(2:2:Nx2, 1:2:Ny2);
ERyy = ER2(1:2:Nx2, 2:2:Ny2);
URzz = UR2(2:2:Nx2, 2:2:Ny2);
```

Calculating the Grid Parameters



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Consideration #1: Wavelength λ_0

The grid resolution must be sufficient to resolve the shortest wavelength.

First, determine the smallest wavelength:

$$\lambda_{\min} = \frac{\min[\lambda_0]}{\max[n(x,y)]}$$

$\min[\lambda_0]$ is the shortest wavelength of interest in the simulation. $\max[n(x,y)]$ is the largest refractive index found anywhere in the grid.

Second, resolve the wave with at least N_λ cells.

$$\Delta_\lambda \approx \frac{\lambda_{\min}}{N_\lambda} \quad N_\lambda \geq 10$$

N_λ	Comments
10 to 20	Low contrast dielectrics
20 to 30	High contrast dielectrics
40 to 60	Most metallic structures
100 to 200	Plasmonic devices

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Consideration #2: Mechanical Features

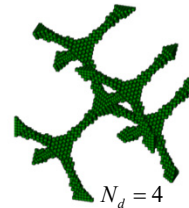
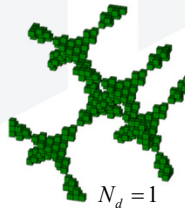
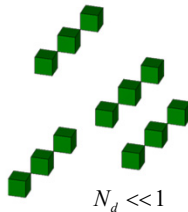
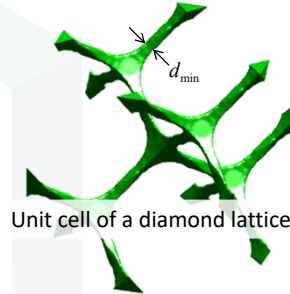
The grid resolution must be sufficient to resolve the smallest mechanical features of the device.

First, determine the smallest feature:

$$d_{\min}$$

Second, resolve d_{\min} by at least N_d cells (usually 1 to 4)

$$\Delta_d \approx \frac{d_{\min}}{N_d} \quad N_d \geq 1$$



Calculating the Initial Grid Resolution

1. Must resolve the minimum wavelength.

$$\Delta_\lambda \leq \frac{\min(\lambda_0) \div \max[n(x, y)]}{N_\lambda} \quad N_\lambda \geq 10$$

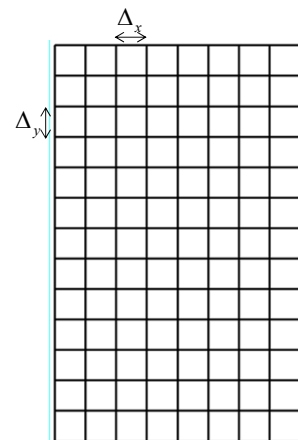
Note: If you are performing a parameter sweep over frequency or wavelength, $\min(\lambda_0)$ is the shortest wavelength in the sweep.

2. Must resolve the minimum structural dimension.

$$\Delta_d \leq \frac{\ell_{\min}}{N_d} \quad N_d \geq 1$$

Proceed with the smallest of the above quantities to be our initial grid resolution

$$\Delta_x = \Delta_y = \min[\Delta_\lambda, \Delta_d]$$

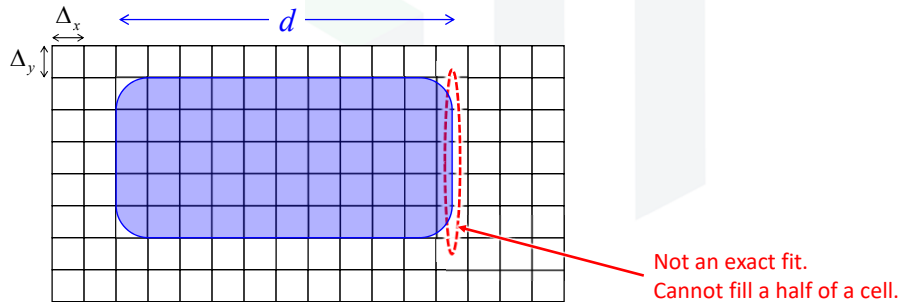


Resolving Critical Dimensions (1 of 3)

The actual dimensions of the device we wish to simulate have not actually yet been considered.

This means we likely cannot resolve the exact dimensions of a device.

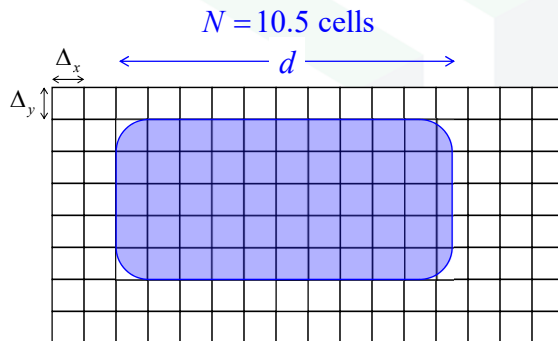
Suppose we wish to place a device of length d onto a grid.



Resolving Critical Dimensions (2 of 3)

To fix this, we first calculate how many cells N are needed to resolve the most important dimension. In this case, let this be d .

$$N = \frac{d}{\Delta_x}$$



Second, we determine how many cells we wish to exactly resolve d . We do this by rounding N up to the nearest integer.

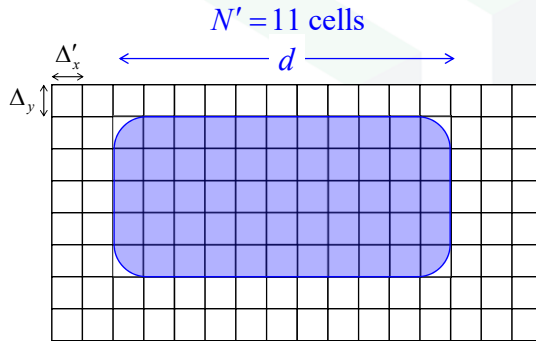
$$N' = \text{ceil} \left[\frac{d}{\Delta_x} \right]$$

$$N' = 11 \text{ cells}$$

Resolving Critical Dimensions (3 of 3)

Third, we adjust the value of Δ_x to represent the dimension d exactly.

$$\Delta'_x = \frac{d}{N'}$$



This step can be called “snapping” the grid to a critical dimension.

Unfortunately, using a uniform grid, it is only possible to do this for one dimension per axis.

“Snap” Grid to Critical Dimensions

Decide what dimensions along each axis are critical.

d_x and d_y

- Typically this is a lattice constant or grating period along x
- Typically this is a film thickness or grating depth along y

Compute how many grid cells comprise d_x and d_y and round UP.

$$M_x = \text{ceil}(d_x/\Delta_x)$$

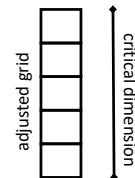
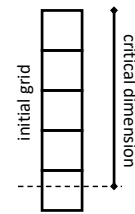
$$M_y = \text{ceil}(d_y/\Delta_y)$$

Adjust grid resolution to fit these critical dimensions in grid EXACTLY.

$$\Delta'_x = d_x/M_x$$

$$\Delta'_y = d_y/M_y$$

For simplicity, drop the prime symbol ' from Δ_x and Δ_y .



Compute Total Grid Size

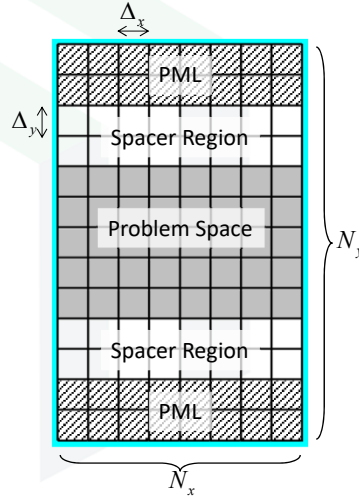
- Don't forget to add cells for PML!
- Must often add "space" between PML and device.

Note: This is particularly important when modeling devices with large evanescent fields.

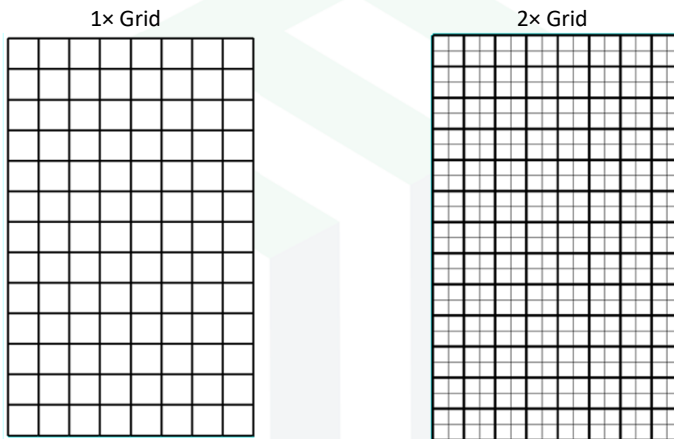
$$N_x = \frac{\Lambda_x}{\Delta_x} \leftarrow \begin{array}{l} \text{Easiest to make } N_x \text{ odd.} \\ \text{Reason discussed later.} \end{array}$$

$$N_y = \frac{\Lambda_y}{\Delta_y} + 2N_{\text{PML}} + 2N_{\text{SPACE}}$$

$$N_{\text{SPACE}} \cong \text{ceil} \left(\frac{\lambda_{\text{max}}}{n_{\text{buff}} \Delta_y} \right)$$



Compute 2x Grid Parameters



$$\begin{array}{ll} N_{x2} = 2 * N_x; & dx2 = dx / 2; \\ N_{y2} = 2 * N_y; & dy2 = dy / 2; \end{array}$$

Averaging At the Edges

Direct

Smoothed

EMPossible Slide 39

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Building Rectangular Structures

For rectangular structures, considering calculating start and stop indices.

Be **very** careful with how many points on the grid you fill in!

```
ny2 = ny1 + round(d/dy2) - 1;
```

Without subtracting 1 here, filling in ny1 to ny2 would include an extra cell. This can introduce error into your results.

```

% BUILD DEVICE
ER2 = er1*ones(Nx2,Ny2);
ER2 (nx1:nx2,ny1:ny2) = er;
ER2 (nx5:nx6,ny1:ny2) = er;
ER2 (:,ny3:ny4) = er;
ER2 (:,ny4+1:Ny2) = er2;
    
```

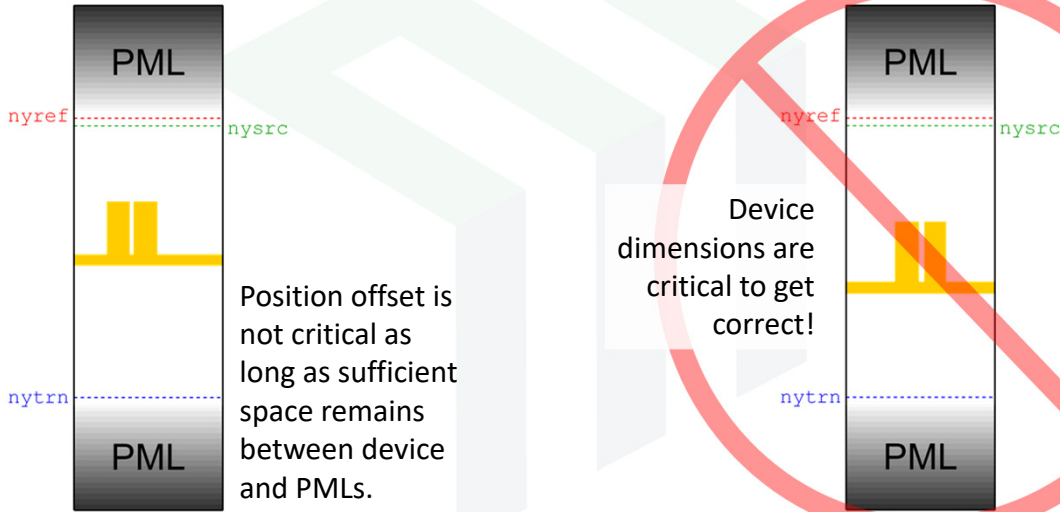
```

%fill everywhere with er1
%add tooth 1
%add tooth 2
%add substrate
%fill transmission region
    
```

EMPossible Slide 40

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Critical and Non-Critical Parameters



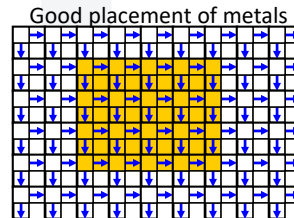
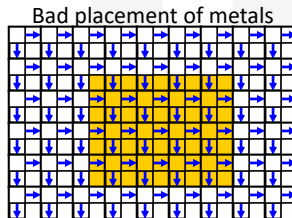
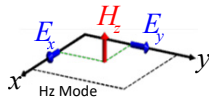
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Oh Yeah, Metals!

Perfect Electric Conductors

$$\epsilon_r = -10000 \quad \text{or} \quad \begin{bmatrix} \ddots & & & & \\ & 0 & 0 & 1 & 0 & 0 \\ & & \ddots & & & \\ & & & E_m & & \\ & & & \vdots & & \\ & & & E_M & & \end{bmatrix} = \begin{bmatrix} f_1 \\ \vdots \\ 0 \\ \vdots \\ f_M \end{bmatrix} \rightarrow E_m = 0$$

Include Tangential Fields at Boundary (TM modes!)



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Walkthrough of the FDFD Algorithm

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Input to the FDFD Algorithm

The FDFD algorithm requires the following information:

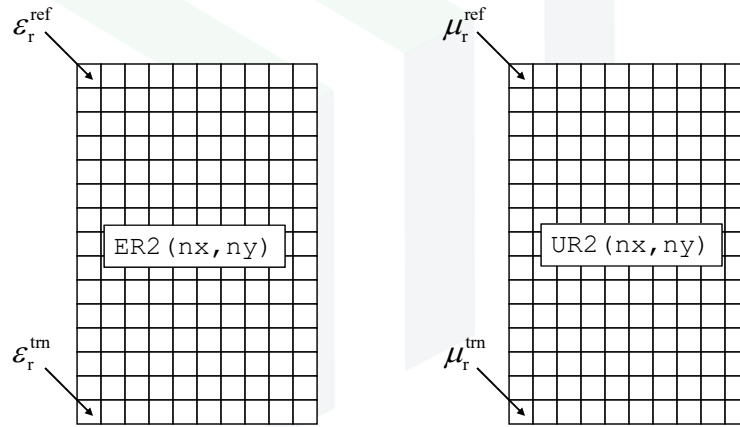
- The materials on the $2\times$ grid: $ER2 (nx, ny)$ and $UR2 (nx, ny)$
- The grid resolution: dx and dy
- The size of the PML on the $1\times$ grid: $NYLO$ and $NYHI$
- The source wavelength, λ_0
- Angle of incidence, θ
- Mode/polarization: 'E' or 'H'

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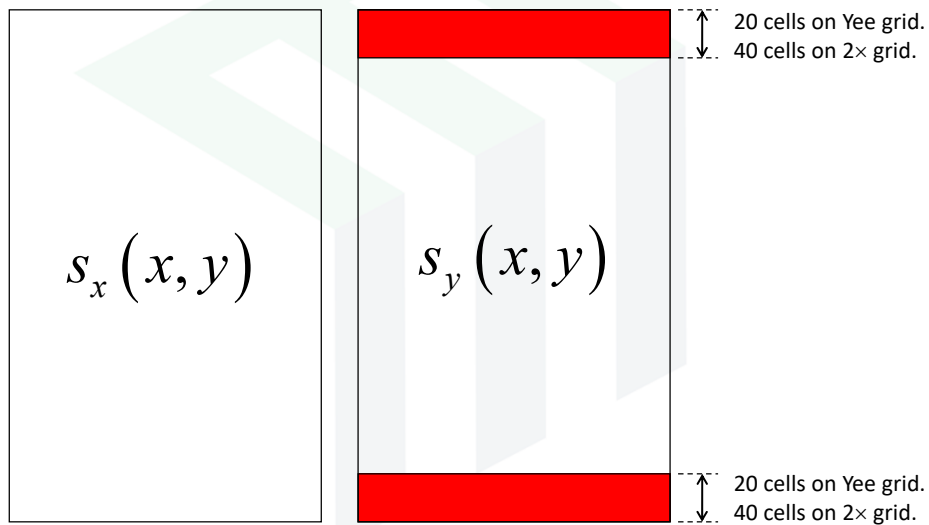
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(1) Determine the Material Properties in the Reflected and Transmitted Regions

If these parameters are not provided in the dashboard, or a dashboard does not exist, they can be pulled directly off of the grid.



(2) Compute the PML Parameters on 2x Grid



(3) Incorporate the PML

We can incorporate the PML parameters into $[\mu]$ and $[\epsilon]$ as follows

$$\nabla \times \vec{E} = k_0 [\mu_r'] \vec{H}$$

$$\nabla \times \vec{H} = k_0 [\epsilon_r'] \vec{E}$$

$$[\epsilon_r'] = \begin{bmatrix} \frac{s_y s_z}{s_x} \epsilon_{xx} & 0 & 0 \\ 0 & \frac{s_x s_z}{s_y} \epsilon_{yy} & 0 \\ 0 & 0 & \frac{s_x s_y}{s_z} \epsilon_{zz} \end{bmatrix}$$

$$[\mu_r'] = \begin{bmatrix} \frac{s_y s_z}{s_x} \mu_{xx} & 0 & 0 \\ 0 & \frac{s_x s_z}{s_y} \mu_{yy} & 0 \\ 0 & 0 & \frac{s_x s_y}{s_z} \mu_{zz} \end{bmatrix}$$

For 2D simulations, $s_z = 1$ and we have

$$\epsilon'_{xx} = \frac{s_y}{s_x} \epsilon_r$$

$$\epsilon'_{yy} = \frac{s_x}{s_y} \epsilon_r$$

$$\epsilon'_{zz} = s_x s_y \epsilon_r$$

$$\mu'_{xx} = \frac{s_y}{s_x} \mu_r$$

$$\mu'_{yy} = \frac{s_x}{s_y} \mu_r$$

$$\mu'_{zz} = s_x s_y \mu_r$$

Note: the PML is incorporated into the $2 \times$ grid.

```

% INCORPORATE PML
URxx = UR2./sx.*sy;
URyy = UR2.*sx./sy;
URzz = UR2.*sx.*sy;
ERxx = ER2./sx.*sy;
ERyy = ER2.*sx./sy;
ERzz = ER2.*sx.*sy;
    
```



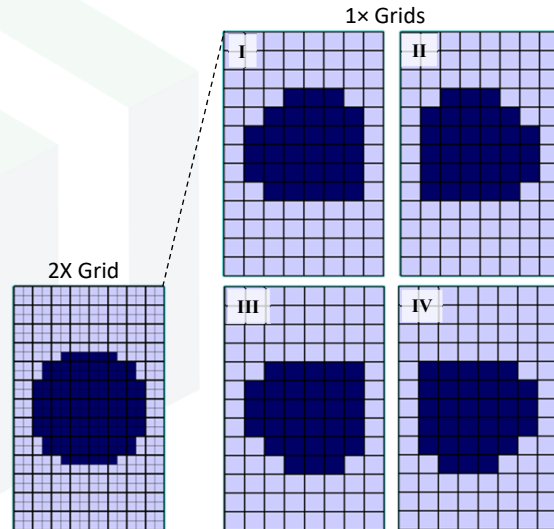
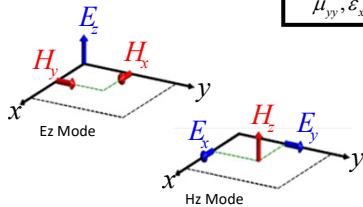
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(4) Overlay Materials Onto 1x Grids

Field and materials assignments

I	II
E_z	H_x, E_y
ϵ_{zz}	μ_{xx}, ϵ_{yy}
III	IV
H_y, E_x	H_z
μ_{yy}, ϵ_{xx}	μ_{zz}



```

% OVERLAY MATERIALS ONTO 1X GRID
URxx = URxx(1:2:Nx2, 2:2:Ny2);
URyy = URyy(2:2:Nx2, 1:2:Ny2);
URzz = URzz(2:2:Nx2, 2:2:Ny2);
ERxx = ERxx(2:2:Nx2, 1:2:Ny2);
ERyy = ERyy(1:2:Nx2, 2:2:Ny2);
ERzz = ERzz(1:2:Nx2, 1:2:Ny2);
    
```



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(5) Compute the Wave Vector Terms

$$k_0 = \frac{2\pi}{\lambda_0}$$

$$\vec{k}_{inc} = k_0 n_{ref} \begin{bmatrix} \sin \theta \\ \cos \theta \end{bmatrix}$$

This is a vector quantity

$$k_x(m) = k_{x,inc} - m \frac{2\pi}{\Lambda_x}$$

$$m = \dots, -2, -1, 0, 1, 2, \dots$$

N_x points total.
For proper symmetry, N_x should be odd.

$$m = [-\text{floor}(N_x/2) : \text{floor}(N_x/2)]';$$

$$k_{y,ref}(m) = \sqrt{(k_0 n_{ref})^2 - k_x^2(m)}$$

$$k_{y,tm}(m) = \sqrt{(k_0 n_{tm})^2 - k_x^2(m)}$$

These equations come from the dispersion equation for the reflected and transmitted regions.

Recall that there used to be a negative sign here. We are able to drop it as long as we also drop the negative sign when calculating diffraction efficiency.

(6) Construct Diagonal Materials Matrices

$$\epsilon_{xx} = \begin{bmatrix} \epsilon_{xx}(1) & & 0 \\ & \epsilon_{xx}(2) & \\ & & \ddots \\ 0 & & & \epsilon_{xx}(N) \end{bmatrix}$$

$$\epsilon_{yy} = \begin{bmatrix} \epsilon_{yy}(1) & & 0 \\ & \epsilon_{yy}(2) & \\ & & \ddots \\ 0 & & & \epsilon_{yy}(N) \end{bmatrix}$$

$$\epsilon_{zz} = \begin{bmatrix} \epsilon_{zz}(1) & & 0 \\ & \epsilon_{zz}(2) & \\ & & \ddots \\ 0 & & & \epsilon_{zz}(N) \end{bmatrix}$$

$$\mu_{xx} = \begin{bmatrix} \mu_{xx}(1) & & 0 \\ & \mu_{xx}(2) & \\ & & \ddots \\ 0 & & & \mu_{xx}(N) \end{bmatrix}$$

$$\mu_{yy} = \begin{bmatrix} \mu_{yy}(1) & & 0 \\ & \mu_{yy}(2) & \\ & & \ddots \\ 0 & & & \mu_{yy}(N) \end{bmatrix}$$

$$\mu_{zz} = \begin{bmatrix} \mu_{zz}(1) & & 0 \\ & \mu_{zz}(2) & \\ & & \ddots \\ 0 & & & \mu_{zz}(N) \end{bmatrix}$$

$$ER_{xx} = \text{diag}(\text{sparse}(ER_{xx}(:)));$$

(7) Construct the Derivative Matrices

$\mathbf{D}_{x'}^e$, $\mathbf{D}_{y'}^e$, $\mathbf{D}_{x'}^h$, and $\mathbf{D}_{y'}^h$

Frequency (or wavelength)
information is incorporated
into FFD here

`[DEX, DEY, DHX, DHY] = yeeder(NGRID, k0*RES, BC, kinc/k0);`

Be sure this function uses sparse matrices from the very beginning.

`A = sparse(M,M)` – creates a sparse M×M matrix of zeros.

`A = spdiags(b,d,A)` – Adds array `b` to diagonal `d` in matrix `A`.

(8) Compute the Wave Matrix A

E Mode

$$\mathbf{A}_E = \mathbf{D}_{x'}^h \boldsymbol{\mu}_{yy}^{-1} \mathbf{D}_{x'}^e + \mathbf{D}_{y'}^h \boldsymbol{\mu}_{xx}^{-1} \mathbf{D}_{y'}^e + \boldsymbol{\epsilon}_{zz}$$

H Mode

$$\mathbf{A}_H = \mathbf{D}_{x'}^e \boldsymbol{\epsilon}_{yy}^{-1} \mathbf{D}_{x'}^h + \mathbf{D}_{y'}^e \boldsymbol{\epsilon}_{xx}^{-1} \mathbf{D}_{y'}^h + \boldsymbol{\mu}_{zz}$$

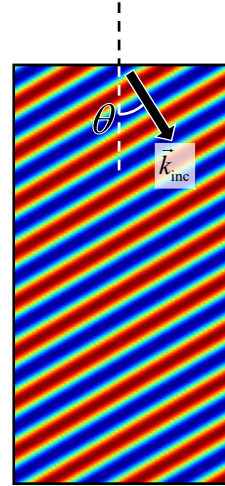
(9) Compute the Source Field

The source has an amplitude of 1.0

$$f_{\text{src}}(x, y) = \exp(j\vec{k}_{\text{inc}} \bullet \vec{r})$$

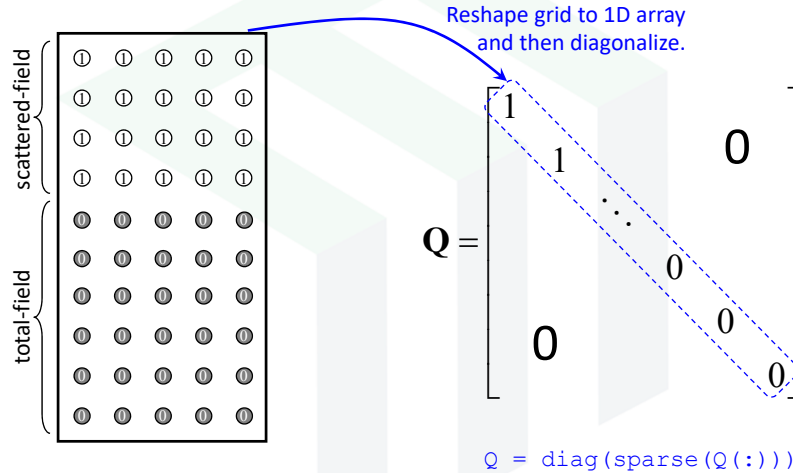
$$= \exp\left[j(k_{x,\text{inc}}x + k_{y,\text{inc}}y)\right]$$

Don't forget to make \mathbf{f}_{src} a column vector.



1x grid

(10) Compute the Scattered-Field Masking Matrix, Q



It is good practice to make the scattered-field region at least one cell larger than the y-low PML.

(11) Compute the source vector, \mathbf{b}

$$\mathbf{A}\mathbf{f} = \mathbf{b} \quad \mathbf{b} = (\mathbf{Q}\mathbf{A} - \mathbf{A}\mathbf{Q})\mathbf{f}_{\text{src}}$$

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(12) Compute the Field \mathbf{f}

$$\mathbf{f} = \mathbf{A}^{-1}\mathbf{b}$$

Aside

In MATLAB, $\mathbf{f} = \mathbf{A} \setminus \mathbf{b}$ employs a direct LU decomposition to calculate \mathbf{f} . This is very stable and robust, but a half-full matrix is created so memory can explode for large problems.

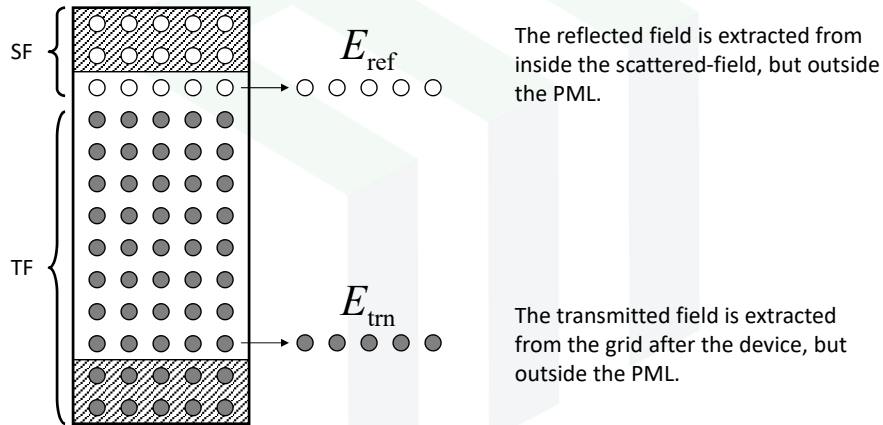
Iterative solutions can be faster and require much less memory, but they are less stable and may never converge to a solution.

Correcting these problems requires significant modification to the FDFD algorithm taught here.

Don't forget to `reshape()` \mathbf{f} from a column vector to a 2D grid after the calculation.

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(13) Extract Transmitted and Reflected Fields



(14) Remove the Phase Tilt

Recall Bloch's theorem,

$$\vec{E}(\vec{r}) = \vec{A}_{k_{inc}}(\vec{r}) \cdot e^{jk_{inc} \cdot \vec{r}}$$

This implies the transmitted and reflected fields have the following form

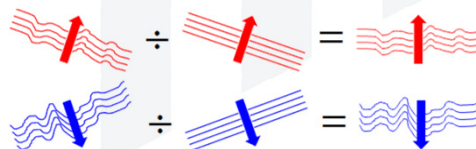
$$E_{ref}(x) = A_{ref}(x) e^{jk_{x,inc} x}$$

$$E_{trn}(x) = A_{trn}(x) e^{jk_{x,inc} x}$$

We remove the phase tilt to calculate the amplitude terms.

$$A_{ref}(x) = E_{ref}(x) e^{-jk_{x,inc} x}$$

$$A_{trn}(x) = E_{trn}(x) e^{-jk_{x,inc} x}$$



(15) Calculate the Complex Amplitudes of the Spatial Harmonics

Recall that the plane wave spectrum is the Fourier transform of the field.



We calculate the FFT of the field amplitude arrays.

$$S_{\text{ref}}(m) = \text{FFT}\{A_{\text{ref}}(x)\}$$

$$S_{\text{tn}}(m) = \text{FFT}\{A_{\text{tn}}(x)\}$$

Some FFT algorithms (like MATLAB) require that you divide by the number of points and shift after calculation.

```
Sref = flipud(fftshift(fft(Aref)))/Nx;
Strn = flipud(fftshift(fft(Atrn)))/Nx;
```

(16) Calculate Diffraction Efficiencies

The source wave was given unit amplitude so

$$|S_{\text{inc}}|^2 = 1$$

The diffraction efficiencies of the reflected modes are then

$$R(m) = |S_{\text{ref}}(m)|^2 \frac{\text{Re}\left[\frac{k_y^{\text{ref}}(m)}{\mu_{r,\text{inc}}}\right]}{\text{Re}\left[\frac{k_y^{\text{inc}}}{\mu_{r,\text{inc}}}\right]} \quad \text{E Mode}$$

$$R(m) = |\tilde{U}_{\text{ref}}(m)|^2 \frac{\text{Re}\left[\frac{k_y^{\text{ref}}(m)}{\varepsilon_{r,\text{inc}}}\right]}{\text{Re}\left[\frac{k_y^{\text{inc}}}{\varepsilon_{r,\text{inc}}}\right]} \quad \text{H Mode}$$

Recall that there used to be a negative sign here. We dropped it because we also dropped the sign when calculating $k_{y,\text{ref}}(m)$.

The diffraction efficiencies of the transmitted modes are then

$$T(m) = |S_{\text{tn}}(m)|^2 \frac{\text{Re}\left[\frac{k_y^{\text{tn}}}{\mu_{r,\text{tn}}}\right]}{\text{Re}\left[\frac{k_y^{\text{inc}}}{\mu_{r,\text{inc}}}\right]} \quad \text{(E Mode)}$$

$$T(m) = |U_{\text{tn}}(m)|^2 \frac{\text{Re}\left[\frac{k_y^{\text{tn}}}{\varepsilon_{r,\text{tn}}}\right]}{\text{Re}\left[\frac{k_y^{\text{inc}}}{\varepsilon_{r,\text{inc}}}\right]} \quad \text{(H Mode)}$$

$S(m)$ \equiv amplitudes of E mode spatial harmonics
 $U(m)$ \equiv amplitudes of H mode spatial harmonics

Note: these equations assume that $|\tilde{S}_{\text{inc}}| = 1$.

(17) Calculate Reflectance, Transmittance, and Conservation of Power

The overall reflectance is

$$REF = \sum_m R(m)$$

The overall transmittance is

$$TRN = \sum_m T(m)$$

Conservation of power is computed as

$$REF + TRN + ABS = 100\%$$

If no loss or gain is incorporated, then $ABS = 0$ and we will have

$$REF + TRN = 100\%$$

- REF + TRN < 100% loss
- REF + TRN = 100% no loss or gain
- REF + TRN > 100% gain

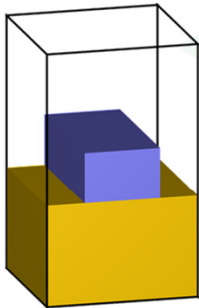


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Remember the Third Dimension!

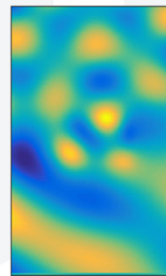
We grabbed a unit cell of a 3D device.



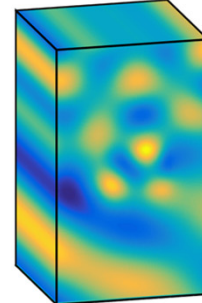
We represented it on a 2D grid.



We simulated it on a 2D grid.



The field is interpreted as infinitely extruded along the third dimension.



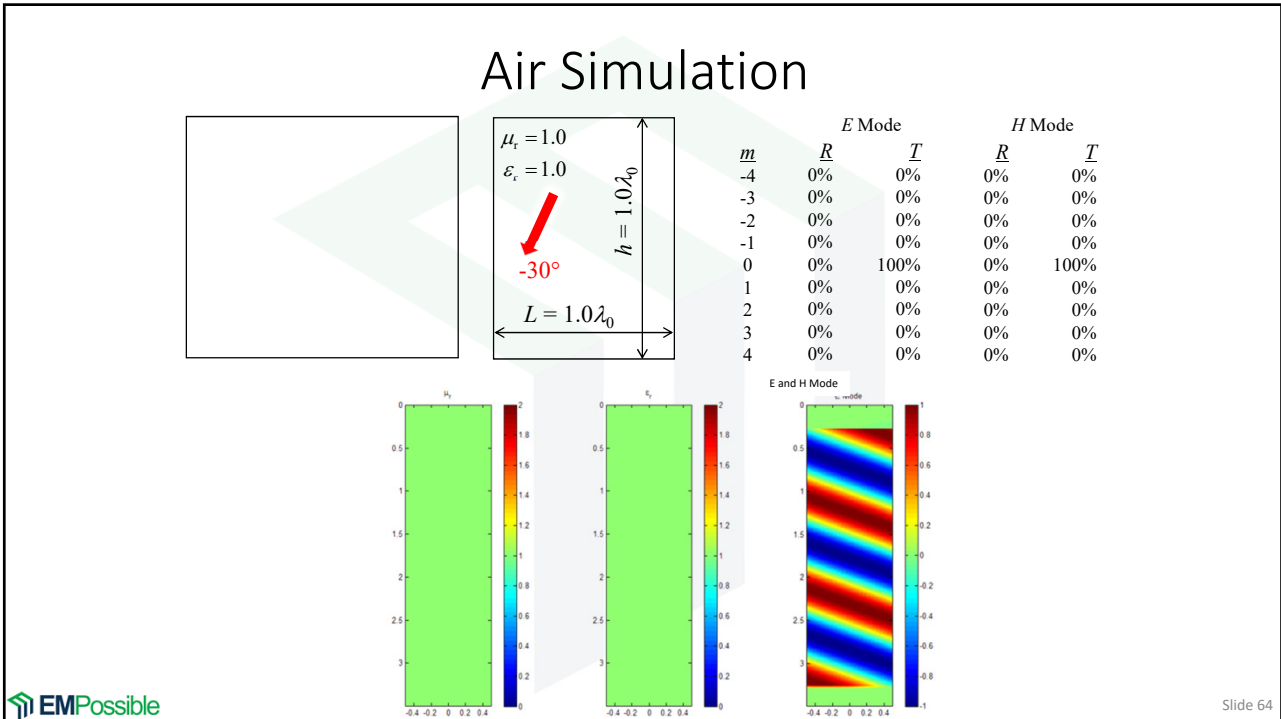
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Examples for Benchmarking

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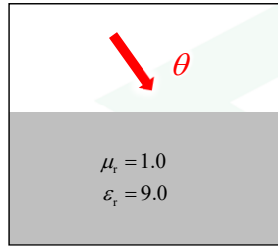
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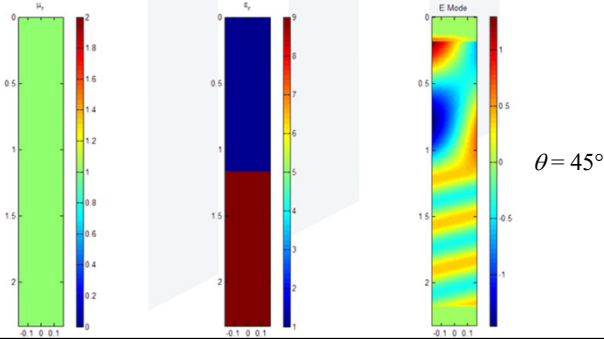
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Dielectric Slab Grating

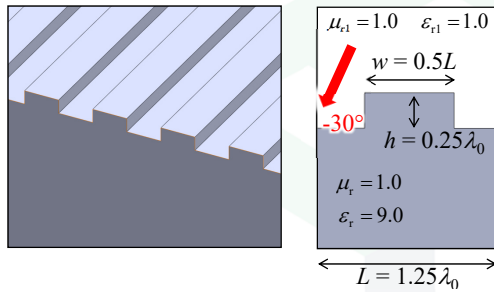


θ	E Mode		H Mode	
	\underline{R}	\underline{T}	\underline{R}	\underline{T}
-60°	49.5%	50.6%	5.0%	95.4%
-25°	28.5%	71.7%	21.8%	78.4%
0°	25.1%	75.0%	25.1%	75.1%
15°	26.3%	73.8%	24.0%	76.3%
45°	37.2%	62.9%	13.9%	86.4%
80°	77.7%	21.8%	8.3%	91.9%

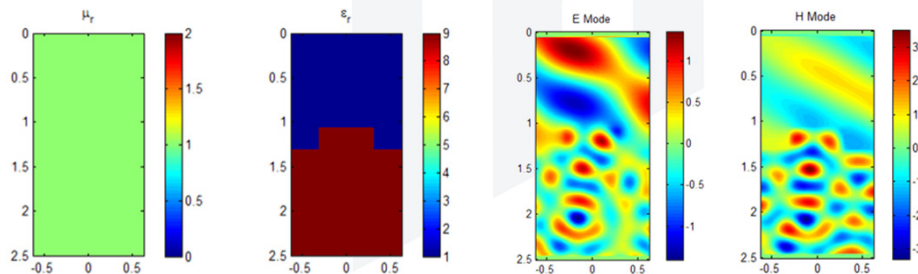
Note: you can come up with your own benchmarking examples using the transfer matrix method!



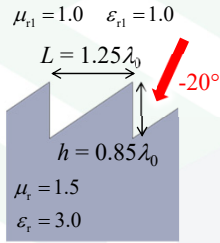
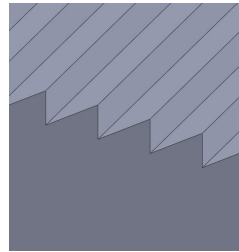
Binary Diffraction Grating



m	E Mode		H Mode	
	\underline{R}	\underline{T}	\underline{R}	\underline{T}
-4	0%	4.6%	0%	2.4%
-3	0%	5.3%	0%	7.5%
-2	0%	2.6%	0%	3.1%
-1	6.9%	22.1%	14.1%	38.8%
0	12.7%	7.8%	0.9%	5.4%
1	0%	17.7%	0%	17.7%
2	0%	16.4%	0%	7.6%
3	0%	3.9%	0%	2.4%
4	0%	0%	0%	0%



Sawtooth Diffraction Grating



m	E Mode		H Mode	
	\underline{R}	\underline{T}	\underline{R}	\underline{T}
-4	0%	0%	0%	0%
-3	0%	0%	0%	0%
-2	0%	5.6%	0%	0.4%
-1	0%	39.2%	0%	53.1%
0	0%	0.8%	0%	3.4%
1	0%	44.2%	0%	29.5%
2	0%	10.1%	0%	13.6%
3	0%	0%	0%	0%
4	0%	0%	0%	0%

