Electromagnetics:
Electromagnetic Field Theory

The Rectangular Waveguide

Lecture Outline

• What is a rectangular waveguide?
• TM Analysis
• TE Analysis
• Visualization of Modes
• Conclusions
What is a Rectangular Waveguide?

Geometry of Rectangular Waveguide

Standard size convention: $a \geq b$
Rectangular waveguides are analyzed a bit like each axis were its own parallel plate waveguide.

Notes on the Rectangular Waveguide

- Most classic waveguide example
- Some of the first waveguides used for microwaves
- Not a transmission line because it has only one conductor
- Does not support a TEM mode
- Exhibits a low-frequency cutoff below which no waves will propagate
Recall TE Analysis

The governing equation for TE analysis is

\[
\frac{\partial^2 H_{0,z}}{\partial x^2} + \frac{\partial^2 H_{0,z}}{\partial y^2} - k_c^2 H_{0,z} = 0 \quad \quad k_c^2 = k^2 - \beta^2
\]

After a solution is obtained, the remaining field components are calculated according to

\[
H_{0,x} = \frac{j\beta}{k_c} \frac{\partial H_{0,z}}{\partial x}, \quad E_{0,x} = \frac{j\omega \mu}{k_c} \frac{\partial H_{0,z}}{\partial y}, \quad E_{0,z} = 0
\]

\[
H_{0,y} = \frac{j\beta}{k_c} \frac{\partial H_{0,z}}{\partial y}, \quad E_{0,y} = \frac{j\omega \mu}{k_c} \frac{\partial H_{0,z}}{\partial x}
\]
General Form of the Solution

From the geometry of the waveguide, the general form of the solution can be immediately written as

\[ H_z(x, y, z) = H_{0,z}(x, y)e^{-j\beta z} \]

Viewing the rectangular waveguide as the combination of two parallel plate waveguides, apply separation of variables to write \( H_{0,z}(x,y) \) as the product of two functions.

\[ H_{0,z}(x, y) = X(x)Y(y) \]

Separation of Variables (1 of 3)

The solution is written as the product of two 1D functions, \( X(x) \) and \( Y(y) \). Substitute this solution back into the differential equation.

\[ \frac{\partial^2 H_{0,z}}{\partial x^2} + \frac{\partial^2 H_{0,z}}{\partial y^2} - k_z^2 H_{0,z} = 0 \]

To be compact, drop the \( (x) \) and \( (y) \) notation.

\[ \frac{\partial^2 XY}{\partial x^2} + \frac{\partial^2 XY}{\partial y^2} - k_z^2 XY = 0 \]

Move \( X(x) \) out of the \( \partial / \partial x \) operation and \( Y(y) \) out of the \( \partial / \partial y \) operation.

\[ \frac{\partial^2 X}{\partial x^2} Y + X \frac{\partial^2 Y}{\partial y^2} - k_z^2 XY = 0 \]

The derivatives become ordinary because \( X(x) \) and \( Y(y) \) have only one independent variable each.
Separation of Variables (2 of 3)

First, attention is focused on the x-dependence in the differential equation.

\[
\frac{1}{X} \frac{d^2X}{dx^2} + \frac{1}{Y} \frac{d^2Y}{dy^2} - k_x^2 = 0
\]

This definition of \( k_x \) lets the differential equation be written as a wave equation.

\[
\frac{d^2X}{dx^2} - k_x^2 X = 0
\]

Second, attention is focused on the y-dependence in the differential equation.

\[
\frac{1}{X} \frac{d^2X}{dx^2} - k_y^2 + \frac{1}{Y} \frac{d^2Y}{dy^2} = 0
\]

This definition of \( k_y \) lets the differential equation be written as a wave equation.

\[
\frac{d^2Y}{dy^2} - k_y^2 Y = 0
\]

Separation of Variables (3 of 3)

If all of this is correct, then it should be possible to add the two new differential equations together to get the original differential equation.

\[
\frac{d^2X}{dx^2} - k_x^2 X = 0 \quad \quad \quad \quad \quad \quad \frac{1}{X} \frac{d^2X}{dx^2} - k_x^2 = 0
\]

\[
\frac{d^2Y}{dy^2} - k_y^2 Y = 0 \quad \quad \quad \quad \quad \quad \frac{1}{Y} \frac{d^2Y}{dy^2} - k_y^2 = 0
\]

\[
\frac{1}{X} \frac{d^2X}{dx^2} + \frac{1}{Y} \frac{d^2Y}{dy^2} - k_x^2 - k_y^2 = 0
\]

The original differential equation is obtained if

\[
k_c^2 = k_x^2 + k_y^2
\]

Original differential equation

\[
\frac{1}{X} \frac{d^2X}{dx^2} + \frac{1}{Y} \frac{d^2Y}{dy^2} - k_c^2 = 0
\]
General Solution

There are now two differential equations to solve.

\[
\frac{d^2X}{dx^2} - k_x^2 X = 0 \quad \text{and} \quad \frac{d^2Y}{dy^2} - k_y^2 Y = 0
\]

These are essentially the same differential equation so their solution has the same general form.

\[
\frac{d^2X}{dx^2} - k_x^2 X = 0 \quad \rightarrow \quad X(x) = A \cos(k_x x) + B \sin(k_x x) \quad \text{PP waveguide along } x
\]

\[
\frac{d^2Y}{dy^2} - k_y^2 Y = 0 \quad \rightarrow \quad Y(y) = C \cos(k_y y) + D \sin(k_y y) \quad \text{PP waveguide along } y
\]

The overall solution is the product of \(X(x)\) and \(Y(y)\).

\[
H_{0,\pm}(x, y) = X(x)Y(y) = [A \cos(k_x x) + B \sin(k_x x)][C \cos(k_y y) + D \sin(k_y y)]
\]

Electromagnetic Boundary Conditions

Boundary conditions require that the tangential component of the electric field be zero at the boundary with a perfect conductor.
In order to apply the boundary conditions, the electric field components $E_{0,x}$ and $E_{0,y}$ must be derived from the expression for $H_{0,z}$.

\[ E_{0,x}(x, y) = \frac{j \omega \mu}{k_e^2} \frac{\partial}{\partial y} H_{0,z} \]
\[ = \frac{j \omega \mu}{k_e^2} \left[ A \cos(k_x x) + B \sin(k_x x) \right] \left[ C \cos(k_y y) + D \sin(k_y y) \right] \]
\[ = -\frac{j \omega \mu}{k_e^2} k_y \left[ A \cos(k_x x) + B \sin(k_x x) \right] \left[ -C \sin(k_y y) + D \cos(k_y y) \right] \]

\[ E_{0,y}(x, y) = \frac{j \omega \mu}{k_e^2} \frac{\partial}{\partial x} H_{0,z} \]
\[ = \frac{j \omega \mu}{k_e^2} \left[ A \cos(k_x x) + B \sin(k_x x) \right] \left[ C \cos(k_y y) + D \sin(k_y y) \right] \]
\[ = \frac{j \omega \mu}{k_e^2} k_x \left[ -A \sin(k_x x) + B \cos(k_x x) \right] \left[ C \cos(k_y y) + D \sin(k_y y) \right] \]

Apply Boundary Conditions (1 of 2)

At the $x = 0$ boundary,

\[ 0 = E_{0,y}(0, y) \]
\[ = \frac{j \omega \mu}{k_e^2} k_x \left[ -A \sin(0) + B \cos(0) \right] \left[ C \cos(k_y y) + D \sin(k_y y) \right] \]
\[ = \frac{j \omega \mu}{k_e^2} k_x \left[ B \left[ C \cos(k_y y) + D \sin(k_y y) \right] \right] \]

\[ B = 0 \]

At the $x = a$ boundary,

\[ 0 = E_{0,y}(a, y) \]
\[ = \frac{j \omega \mu}{k_e^2} k_x \left[ -A \sin(k_x a) \right] \left[ C \cos(k_y y) + D \sin(k_y y) \right] \]

$A = 0$ leads to a trivial solution. It must be the $\sin(k_x a)$ term that enforces the BC.

\[ 0 = \sin(k_x a) \quad \rightarrow \quad k_x a = m \pi \quad m = 0, 1, 2, ... \]

The specific values $m$ can be will be considered later.
Apply Boundary Conditions (2 of 2)

At the \( y = 0 \) boundary,

\[
0 = E_{0,x}(x,0) = \frac{j \omega}{k_c} k_x \left[ A \cos(k_x x) + B \sin(k_x x) \right] \left[ -C \sin(0) + D \cos(0) \right] = -\frac{j \omega}{k_c} k_x \left[ A \cos(k_x x) + B \sin(k_x x) \right] \left[ D \right] \quad \Rightarrow \quad D = 0
\]

At the \( y = b \) boundary,

\[
0 = E_{0,y}(x,b) = -\frac{j \omega}{k_c} k_y \left[ A \cos(k_x x) + B \sin(k_x x) \right] \left[ -C \sin(k_y b) \right]
\]

\( C = 0 \) leads to a trivial solution. It must be the \( \sin(k_y b) \) term that enforces the BC.

\[
0 = \sin(k_y b) \quad \Rightarrow \quad k_y b = n \pi \quad n = 0, 1, 2, \ldots
\]

The specific values \( n \) can be will be considered later.

Revised Solution for \( H_{0,z} \)

It was determined that \( B = D = 0 \) so the expression for \( H_{0,z} \) becomes

\[
H_{0,z}(x, y) = \frac{AC}{k_c} \cos(k_x x) \cos(k_y y)
\]

The product \( AC \) is written as a single constant \( A_{mn} \).

\[
H_{0,z}(x, y) = A_{mn} \cos(k_x x) \cos(k_y y)
\]

Also, recall the conditions for \( k_x \) and \( k_y \).

\[
k_x a = m \pi \quad \Rightarrow \quad k_x = \frac{m \pi}{a} \\
k_y b = n \pi \quad \Rightarrow \quad k_y = \frac{n \pi}{b}
\]

\[
H_{0,z}(x, y) = A_{mn} \cos \left( \frac{m \pi x}{a} \right) \cos \left( \frac{n \pi y}{b} \right)
\]
Entire Solution (1 of 2)

The final expression for $H_{0,z}$ is

$$H_{0,z}(x, y) = A_{mn} \cos \left( \frac{m \pi x}{a} \right) \cos \left( \frac{n \pi y}{b} \right)$$

$E_{0,z}(x, y) = 0$

From this, the other field components are

$$E_{0,x}(x, y) = \frac{\omega \mu m \pi}{k \pi} A_{mn} \cos \left( \frac{m \pi x}{a} \right) \sin \left( \frac{n \pi y}{b} \right)$$

$$E_{0,y}(x, y) = -\frac{\omega \mu m \pi}{k \pi} A_{mn} \sin \left( \frac{m \pi x}{a} \right) \cos \left( \frac{n \pi y}{b} \right)$$

$$H_{0,x}(x, y) = \frac{j \beta m \pi}{k \pi} A_{mn} \sin \left( \frac{m \pi x}{a} \right) \cos \left( \frac{n \pi y}{b} \right)$$

$$H_{0,y}(x, y) = \frac{j \beta m \pi}{k \pi} A_{mn} \cos \left( \frac{m \pi x}{a} \right) \sin \left( \frac{n \pi y}{b} \right)$$

Entire Solution (2 of 2)

The overall electric and magnetic fields at any position are

$$E_x(x, y, z) = \frac{j \omega \mu m \pi}{k \pi} A_{mn} \cos \left( \frac{m \pi x}{a} \right) \sin \left( \frac{n \pi y}{b} \right) e^{-j \beta \omega z}$$

$$E_y(x, y, z) = -\frac{j \omega \mu m \pi}{k \pi} A_{mn} \sin \left( \frac{m \pi x}{a} \right) \cos \left( \frac{n \pi y}{b} \right) e^{-j \beta \omega z}$$

$$E_z(x, y, z) = 0$$

$$H_x(x, y, z) = \frac{j \beta m \pi}{k \pi} A_{mn} \sin \left( \frac{m \pi x}{a} \right) \cos \left( \frac{n \pi y}{b} \right) e^{-j \beta \omega z}$$

$$H_y(x, y, z) = \frac{j \beta m \pi}{k \pi} A_{mn} \cos \left( \frac{m \pi x}{a} \right) \sin \left( \frac{n \pi y}{b} \right) e^{-j \beta \omega z}$$

$$H_z(x, y, z) = A_{mn} \cos \left( \frac{m \pi x}{a} \right) \cos \left( \frac{n \pi y}{b} \right) e^{-j \beta \omega z}$$
Phase Constant, $\beta$

Recall the cutoff wave number

$$k_c^2 = k_x^2 + k_y^2$$

After analyzing the boundary conditions, this expression can be written as

$$k_c^2 = \left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2$$

The phase constant $\beta$ is therefore

$$k_c^2 = k^2 - \beta^2$$
$$\beta^2 = k^2 - k_c^2$$

$$\beta_{mn} = \sqrt{k^2 - k_c^2} = \sqrt{k^2 - \left(\frac{m\pi}{a}\right)^2 - \left(\frac{n\pi}{b}\right)^2}$$

Cutoff Frequency, $f_c$

Recall the expression for the phase constant

$$\beta_{mn} = \sqrt{k^2 - k_c^2}$$

The phase constant must be a real number for a guided mode. This requires

$$k > k_c$$

Any time $k < k_c$, the mode is cutoff and not supported by the waveguide. From this, the cutoff frequency $f_c$ is derived to be

$$f_{c,mn} = \frac{k_c}{2\pi \sqrt{\mu\varepsilon}} = \frac{1}{2\pi \sqrt{\mu\varepsilon}} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2}$$
Characteristic Impedance, $Z_{TE}$

The characteristic impedance $Z_{TE}$ of the TE mode is

$$Z_{TE} = \frac{E_x}{H_y} = \frac{j\omega \mu \pi}{k_c^2 b} A_{mn} \cos \left( \frac{m\pi x}{a} \right) \sin \left( \frac{n\pi y}{b} \right) e^{-j\beta z} = \frac{\omega \mu}{\beta_{mn}} = \frac{k\eta}{\beta_{mn}}$$

Cutoff for First-Order TE Mode (1 of 2)

The cutoff frequency for the $TE_{mn}$ mode was found to be

$$f_{c, mn} = \frac{1}{2\pi\sqrt{\mu\epsilon}} \sqrt{\left( \frac{m\pi}{a} \right)^2 + \left( \frac{n\pi}{b} \right)^2}$$

What about the $TE_{00}$ mode?

$TE_{00} \rightarrow m = n = 0$

$$f_{c, 00} = \frac{1}{2\pi\sqrt{\mu\epsilon}} \sqrt{\left( \frac{0}{a} \right)^2 + \left( \frac{0}{b} \right)^2} = 0$$

The $TE_{00}$ mode does not exist because it is impossible for a mode in this waveguide to have a cutoff frequency of 0 Hz.
Cutoff for First-Order TE Mode (2 of 2)

What about the TE_{01} mode?

\[ \text{TE}_{01} \rightarrow m = 0, n = 1 \]

\[ f_{c,01} = \frac{1}{2\pi\sqrt{\mu\varepsilon}} \sqrt{\left( \frac{0}{a} \right)^2 + \left( \frac{1 \cdot \pi}{b} \right)^2} = \frac{1}{2b\sqrt{\mu\varepsilon}} \]

What about the TE_{10} mode?

\[ \text{TE}_{10} \rightarrow m = 1, n = 0 \]

\[ f_{c,10} = \frac{1}{2\pi\sqrt{\mu\varepsilon}} \sqrt{\left( \frac{1 \cdot \pi}{a} \right)^2 + \left( \frac{0}{b} \right)^2} = \frac{1}{2a\sqrt{\mu\varepsilon}} \]

Since \( a > b \), it is concluded that the first-order TE mode is TE_{10} because it has the lowest cutoff frequency.

CAUTION: It cannot yet be said that the TE_{10} is the fundamental mode because the cutoff frequency of the TM modes has not yet been checked.

Single Mode Operation (1 of 2)

Over what range of frequencies does a rectangular waveguide supports only a single TE mode?

\( f_{c1} < f < f_{c2} \)

**Low-Frequency Cutoff**

The lower frequency cutoff was just found.

\[ f_{c1} = \frac{1}{2a\sqrt{\mu\varepsilon}} \]

**High-Frequency Cutoff**

The high-frequency cutoff is the frequency where the second-order TE mode is supported. This could be the TE_{01}, TE_{11} or TE_{20} mode. All must be considered.

\[ \begin{align*}
\text{TE}_{01} : f_{c,01} &= \frac{1}{2\pi\sqrt{\mu\varepsilon}} \sqrt{\left( \frac{0}{a} \right)^2 + \left( \frac{1 \cdot \pi}{b} \right)^2} = \frac{1}{2b\sqrt{\mu\varepsilon}} \\
\text{TE}_{11} : f_{c,11} &= \frac{1}{2\pi\sqrt{\mu\varepsilon}} \sqrt{\left( \frac{\pi}{a} \right)^2 + \left( \frac{\pi}{b} \right)^2} = \frac{1}{2b\sqrt{\mu\varepsilon}} \sqrt{1 + \left( \frac{b}{a} \right)^2} \\
\text{TE}_{20} : f_{c,20} &= \frac{1}{2\pi\sqrt{\mu\varepsilon}} \sqrt{\left( \frac{2 \cdot \pi}{a} \right)^2 + \left( \frac{0}{b} \right)^2} = \frac{1}{2b\sqrt{\mu\varepsilon}} \left( \frac{2b}{a} \right) \\
\end{align*} \]

\[ f_{c2} = \begin{cases} f_{c,01} & a \leq 2b \\ f_{c,20} & a > 2b \quad \text{(typical)} \end{cases} \]

TE_{11} will always have a higher cutoff frequency than TE_{01}.

The second-order mode depends on choice of \( a \) and \( b \).
Single Mode Operation (2 of 2)

Bandwidth

Typical rectangular waveguides will have \( a > 2b \), so

\[
\begin{align*}
f_{c1} &= \frac{1}{2a\sqrt{\mu\varepsilon}} \\
f_{c2} &= \frac{1}{a\sqrt{\mu\varepsilon}}
\end{align*}
\]

\[
\Delta f = f_{c2} - f_{c1} = \frac{1}{a\sqrt{\mu\varepsilon}} - \frac{1}{2a\sqrt{\mu\varepsilon}} = \frac{1}{2a\sqrt{\mu\varepsilon}}
\]

Fractional Bandwidth

Continuing the assumption that \( a > 2b \), the fractional bandwidth can be calculated from \( f_{c1} \) and \( f_{c2} \) above as follows

\[
FBW = \frac{\Delta f}{f_c} = \frac{f_{c2} - f_{c1}}{(f_{c2} + f_{c1})/2} = 2 \frac{f_{c2} - f_{c1}}{f_{c2} + f_{c1}} = 2 \frac{1}{a\sqrt{\mu\varepsilon}} - \frac{1}{2a\sqrt{\mu\varepsilon}} = \frac{2}{3} = 66.7\%
\]

Example #1 – TE Mode Analysis (1 of 4)

Suppose there exists an air-filled rectangular waveguide with \( a = 3 \text{ cm} \) and \( b = 2 \text{ cm} \).

What is the cutoff frequency of the waveguide?

\[
f_{c1} = \frac{1}{2a\sqrt{\mu\varepsilon}} = \frac{c_0}{2a\sqrt{\mu\varepsilon}} = \frac{299792458 \text{ m/s}}{2(0.03 \text{ m})\sqrt{(1.0)(1.0)}} = 5.0 \text{ GHz}
\]

Over what range of frequencies is the waveguide single mode?

Observing that \( a < 2b \), so the second-order mode is \( \text{TE}_{01} \).

\[
f_{c2} = \frac{1}{2b\sqrt{\mu\varepsilon}} = \frac{c_0}{2b\sqrt{\mu\varepsilon}} = \frac{299792458 \text{ m/s}}{2(0.02 \text{ m})\sqrt{(1.0)(1.0)}} = 7.5 \text{ GHz}
\]

\[5.0 \text{ GHz} < f < 7.5 \text{ GHz}\]
Example #1 – TE Mode Analysis (2 of 4)

What is the fractional bandwidth of the waveguide?

FBW = \frac{\Delta f}{f} = \frac{(100\%) \frac{f_2 - f_1}{f_2 + f_1}}{(200\%) \frac{7.5 - 5.0}{7.5 + 5.0}} = \frac{40\%}{40\%} = 40\%

Plot the phase constant and effective refractive index for the first-order and second-order modes from DC up to 15 GHz.

The phase constant is calculated as:

\beta_{mn} = \sqrt{k^2 - \left(\frac{m\pi}{a}\right)^2 - \left(\frac{n\pi}{b}\right)^2}

\rightarrow

TE_{01}: \beta_1 = \sqrt{\frac{2\pi f}{c_0} - \left(\frac{\pi}{a}\right)^2}

TE_{02}: \beta_2 = \sqrt{\frac{2\pi f}{c_0} - \left(\frac{\pi}{b}\right)^2}

The effective refractive index is calculated as:

\beta = k_0 n_{\text{eff}} \rightarrow n_{\text{eff}} = \frac{\beta}{\frac{2\pi f}{c_0}} = n_0 \frac{c_0}{2\pi f}

Example #1 – TE Mode Analysis (3 of 4)

Plot the phase constant and effective refractive index for the first-order and second-order modes from DC up to 15 GHz.

\beta_1 = \sqrt{\frac{2\pi f}{c_0} - \left(\frac{\pi}{a}\right)^2}

\beta_2 = \sqrt{\frac{2\pi f}{c_0} - \left(\frac{\pi}{b}\right)^2}

n_{\text{eff,1}} = n_0 \frac{c_0}{2\pi f}

n_{\text{eff,2}} = n_0 \frac{c_0}{2\pi f}
Example #1 – TE Mode Analysis (4 of 4)

Plot the velocity of the modes as a function of frequency.

\[ v_1 = \frac{c_0}{n_{e1}} \quad v_2 = \frac{c_0}{n_{e2}} \]

Are the modes travelling faster than the speed of light?

Summary of TE Analysis

Field Solution
\[
E_x(x, y, z) = \frac{j \omega \mu_0 \pi}{k^2} A_{m}\cos \left( \frac{m\pi x}{a} \right) \sin \left( \frac{n\pi y}{b} \right) e^{-j\beta_z z}
\]
\[
E_y(x, y, z) = \frac{j \omega \mu_0 \pi}{k^2} A_{m}\sin \left( \frac{m\pi x}{a} \right) \cos \left( \frac{n\pi y}{b} \right) e^{-j\beta_z z}
\]
\[
E_z(x, y, z) = 0
\]
\[
H_x(x, y, z) = \frac{j \beta_{m} m\pi}{k^2} A_{m}\cos \left( \frac{m\pi x}{a} \right) \sin \left( \frac{n\pi y}{b} \right) e^{-j\beta_z z}
\]
\[
H_y(x, y, z) = \frac{j \beta_{m} n\pi}{k^2} A_{m}\cos \left( \frac{m\pi x}{a} \right) \sin \left( \frac{n\pi y}{b} \right) e^{-j\beta_z z}
\]
\[
H_z(x, y, z) = A_{m}\cos \left( \frac{m\pi x}{a} \right) \cos \left( \frac{n\pi y}{b} \right) e^{-j\beta_z z}
\]

Phase Constant
\[
\beta_m = \sqrt{k^2 - \left( \frac{m\pi}{a} \right)^2 - \left( \frac{n\pi}{b} \right)^2}
\]
Same equation as for TM

Cutoff Frequency
\[
f_{\text{cut}} = \frac{1}{2\pi \sqrt{\mu \epsilon}} \sqrt{\left( \frac{m\pi}{a} \right)^2 + \left( \frac{n\pi}{b} \right)^2}
\]
Same equation as for TM

Characteristic Impedance
\[
Z_{\text{TE,m}} = \frac{k \omega}{\epsilon \mu \beta_m}
\]
The governing equation for TM analysis is
\[
\frac{\partial^2 E_{0,z}}{\partial x^2} + \frac{\partial^2 E_{0,z}}{\partial y^2} - k_z^2 E_{0,z} = 0 \quad H_{0,z} = 0 \quad k_z^2 = k^2 - \beta^2
\]

After a solution is obtained, the remaining field components are calculated according to
\[
\begin{align*}
H_{0,x} &= \frac{j \omega \epsilon}{k_z} \frac{\partial E_{0,z}}{\partial y} \\
H_{0,y} &= -\frac{j \omega \epsilon}{k_z} \frac{\partial E_{0,z}}{\partial x} \\
E_{0,x} &= -\frac{j \beta}{k_z} \frac{\partial E_{0,z}}{\partial z} \\
E_{0,y} &= -\frac{j \beta}{k_z} \frac{\partial E_{0,z}}{\partial z}
\end{align*}
\]
General Form of the Solution

From the geometry of the waveguide, the general form of the solution can be immediately written as

\[ E_z(x, y, z) = E_{0,z}(x, y)e^{-j\beta z} \]

Viewing the rectangular waveguide as the combination of two parallel plate waveguides, apply separation of variables to write \( E_{0,z}(x, y) \) as the product of two functions.

\[ E_{0,z}(x, y) = X(x)Y(y) \]

Separation of Variables (1 of 3)

The solution is written as the product of two 1D functions, \( X(x) \) and \( Y(y) \). Substitute this solution back into the differential equation.

\[ \frac{\partial^2 E_{0,z}}{\partial x^2} + \frac{\partial^2 E_{0,z}}{\partial y^2} - k_e^2 E_{0,z} = 0 \]

\[ \frac{\partial^2 XY}{\partial x^2} + \frac{\partial^2 XY}{\partial y^2} - k_e^2 XY = 0 \]

To be compact, drop the \( x \) and \( y \) notation.

\[ \frac{\partial^2 X}{\partial x^2}Y + X \frac{\partial^2 Y}{\partial y^2} - k_e^2 XY = 0 \]

Move \( X(x) \) out of the \( \partial / \partial x \) operation and \( Y(y) \) out of the \( \partial / \partial y \) operation.

\[ \frac{1}{X} \frac{d^2 X}{dx^2} + \frac{1}{Y} \frac{d^2 Y}{dy^2} - k_e^2 = 0 \]

The derivatives become ordinary because \( X(x) \) and \( Y(y) \) have only one independent variable each.
Separation of Variables (2 of 3)

First, attention is focused on the $x$-dependence in the differential equation.

\[
\frac{1}{X} \frac{d^2 X}{dx^2} + \frac{1}{Y} \frac{d^2 Y}{dy^2} - k_x^2 = 0
\]

This definition of $k_x$ lets the differential equation be written as a wave equation.

\[
\frac{d^2 X}{dx^2} - k_x^2 X = 0
\]

Second, attention is focused on the $y$-dependence in the differential equation.

\[
\frac{1}{X} \frac{d^2 X}{dx^2} - k_x^2 + \frac{d^2 Y}{Y} \frac{d^2 Y}{dy^2} = 0
\]

This definition of $k_y$ lets the differential equation be written as a wave equation.

\[
\frac{d^2 Y}{dy^2} - k_y^2 Y = 0
\]

Separation of Variables (3 of 3)

It should be possible to add the two new differential equations together to get the original differential equation.

\[
\frac{d^2 X}{dx^2} - k_x^2 X = 0
\]

\[
\frac{d^2 Y}{dy^2} - k_y^2 Y = 0
\]

\[
\frac{1}{X} \frac{d^2 X}{dx^2} + \frac{1}{Y} \frac{d^2 Y}{dy^2} - k_x^2 - k_y^2 = 0
\]

The original differential equation is obtained if

\[
k_x^2 = k_x^2 + k_y^2
\]

Original differential equation

\[
\frac{1}{X} \frac{d^2 X}{dx^2} + \frac{1}{Y} \frac{d^2 Y}{dy^2} - k_x^2 = 0
\]
General Solution

There are now two differential equations to solve.

\[ \frac{d^2X}{dx^2} - k_x^2 X = 0 \quad \frac{d^2Y}{dy^2} - k_y^2 Y = 0 \]

These are essentially the same differential equation so their solution has the same general form.

\[ \frac{d^2X}{dx^2} - k_x^2 X = 0 \quad \rightarrow \quad X(x) = A \cos(k_x x) + B \sin(k_x x) \quad \text{PP waveguide along } x \]

\[ \frac{d^2Y}{dy^2} - k_y^2 Y = 0 \quad \rightarrow \quad Y(y) = C \cos(k_y y) + D \sin(k_y y) \quad \text{PP waveguide along } y \]

The overall solution is the product of \(X(x)\) and \(Y(y)\).

\[ E_{0,z}(x, y) = X(x)Y(y) = [A \cos(k_x x) + B \sin(k_x x)] [C \cos(k_y y) + D \sin(k_y y)] \]

Electromagnetic Boundary Conditions

BCs require that the tangential component of the electric field be zero at the boundary with a perfect conductor. \(E_{0,z}\) is tangential to all interfaces so it is just used directly. There is no need to calculate \(E_{0,x}\) or \(E_{0,y}\).
Apply Boundary Conditions (1 of 2)

At the $x = 0$ boundary,

$$0 = E_{a,x}(0,y)$$

$$= [A \cos(0) + B \sin(0)][C \cos(k_y) + D \sin(k_y)]$$

$$= [A][C \cos(k_y) + D \sin(k_y)]$$

$A = 0$

At the $x = a$ boundary,

$$0 = E_{a,x}(a,y)$$

$$= [B \sin(k_a)][C \cos(k_y) + D \sin(k_y)]$$

$B = 0$ leads to a trivial solution. It must be the $\sin(k_a)$ term that enforces the BC.

$$0 = \sin(k_a) \rightarrow k_a = m\pi \quad m = 1, 2, \ldots \quad m = 0$ leads to a trivial solution.$$

Apply Boundary Conditions (2 of 2)

At the $y = 0$ boundary,

$$0 = E_{a,y}(x,0)$$

$$= [A \cos(k_x) + B \sin(k_x)][C \cos(0) + D \sin(0)]$$

$$= [A \cos(k_x) + B \sin(k_x)][C]$$

$C = 0$

At the $y = b$ boundary,

$$0 = E_{a,y}(x,b)$$

$$= [A \cos(k_x) + B \sin(k_x)][D \sin(k_y)]$$

$D = 0$ leads to a trivial solution. It must be the $\sin(k_y)$ term that enforces the BC.

$$0 = \sin(k_y) \rightarrow k_y = n\pi \quad n = 1, 2, \ldots \quad n = 0$ leads to a trivial solution.$
Revised Solution for $E_{0,z}$

It was determined that $A = C = 0$ so the expression for $E_{0,z}$ becomes

$$E_{0,z}(x, y) = BD \sin(k_x x) \sin(k_y y)$$

The product $BD$ is written as a single constant $B_{mn}$.

$$E_{0,z}(x, y) = B_{mn} \sin(k_x x) \sin(k_y y)$$

Also, recall the conditions for $k_x$ and $k_y$:

$$k_x = \frac{m \pi}{a}, \quad k_y = \frac{n \pi}{b}$$

$$E_{0,z}(x, y) = B_{mn} \sin\left(\frac{m \pi x}{a}\right) \sin\left(\frac{n \pi y}{b}\right)$$

* Note: Neither $m$ nor $n$ can be zero or the entire solution will be zero.

Entire Solution (1 of 2)

The final expression for $E_{0,z}$ is

$$E_{0,z}(x, y) = B_{mn} \sin\left(\frac{m \pi x}{a}\right) \sin\left(\frac{n \pi y}{b}\right) \quad H_{0,z}(x, y) = 0$$

From this, the other field components are

$$E_{0,x}(x, y) = -\frac{j \beta_{mn} m \pi}{k_y^2 a} B_{mn} \cos\left(\frac{m \pi x}{a}\right) \sin\left(\frac{n \pi y}{b}\right)$$

$$E_{0,y}(x, y) = -\frac{j \beta_{mn} n \pi}{k_x^2 b} B_{mn} \sin\left(\frac{m \pi x}{a}\right) \cos\left(\frac{n \pi y}{b}\right)$$

$$H_{0,x}(x, y) = -\frac{j \omega \varepsilon_{mn}}{k_y^2 b} B_{mn} \sin\left(\frac{m \pi x}{a}\right) \cos\left(\frac{n \pi y}{b}\right)$$

$$H_{0,y}(x, y) = -\frac{j \omega \mu_{mn}}{k_x^2 a} B_{mn} \cos\left(\frac{m \pi x}{a}\right) \sin\left(\frac{n \pi y}{b}\right)$$
The overall electric and magnetic fields at any position are

\[ E_x(x, y, z) = -j \beta_{mn} \frac{m \pi}{k_z} B_{mn} \cos\left(\frac{m \pi x}{a}\right) \sin\left(\frac{n \pi y}{b}\right) e^{-j \omega z} \]

\[ E_y(x, y, z) = -j \beta_{mn} \frac{n \pi}{k_z} B_{mn} \sin\left(\frac{m \pi x}{a}\right) \cos\left(\frac{n \pi y}{b}\right) e^{-j \omega z} \]

\[ E_z(x, y, z) = B_{mn} \sin\left(\frac{m \pi x}{a}\right) \sin\left(\frac{n \pi y}{b}\right) e^{-j \omega z} \]

\[ H_x(x, y, z) = \frac{j \omega \epsilon \pi}{k_z^2} \sin\left(\frac{m \pi x}{a}\right) \cos\left(\frac{n \pi y}{b}\right) e^{-j \omega z} \]

\[ H_y(x, y, z) = -\frac{j \omega \epsilon \pi}{k_z^2} \cos\left(\frac{m \pi x}{a}\right) \sin\left(\frac{n \pi y}{b}\right) e^{-j \omega z} \]

\[ H_z(x, y, z) = 0 \]

**Phase Constant, \( \beta \)**

Recall the cutoff wave number

\[ k_z^2 = k_x^2 + k_y^2 \]

After analyzing the boundary conditions, this expression can be written as

\[ k_z^2 = \left(\frac{m \pi}{a}\right)^2 + \left(\frac{n \pi}{b}\right)^2 \]

The phase constant \( \beta \) is therefore

\[ k_z^2 = k^2 - \beta^2 \]

\[ \beta^2 = k_z^2 - k_x^2 \]

\[ \beta_{mn} = \sqrt{k^2 - k_x^2} = \sqrt{k^2 - \left(\frac{m \pi}{a}\right)^2 - \left(\frac{n \pi}{b}\right)^2} \]
Cutoff Frequency, $f_c$

Recall the expression for the phase constant

$$\beta_{mn} = \sqrt{k^2 - k_c^2}$$

The phase constant must be a real number for a guided mode. This requires

$$k > k_c$$

Any time $k < k_c$, the mode is cutoff and not supported by the waveguide. From this, the cutoff frequency $f_c$ can be derived as

$$f_c = \omega / 2\pi = 1 / \omega \sqrt{\mu \varepsilon} = 1 / \omega \sqrt{\mu \varepsilon} \sqrt{m \pi / a + n \pi / b}^2$$

This is the same equation as for the TE modes.

Characteristic Impedance, $Z_{TM}$

The characteristic impedance $Z_{TM}$ for the TM mode is

$$Z_{TM} = \frac{E_x}{H_y} = \frac{j \beta_{mn} m \pi}{k_c^2 a} B_{mn} \cos \left( \frac{m \pi x}{a} \right) \sin \left( \frac{n \pi y}{b} \right) e^{-j \beta z}$$

$$= \frac{j \omega \varepsilon \mu \pi}{k^2 a} B_{mn} \cos \left( \frac{m \pi x}{a} \right) \sin \left( \frac{n \pi y}{b} \right) e^{-j \beta z}$$

$$= \frac{\beta_{mn} \eta}{\omega \varepsilon k}$$
Cutoff for First-Order TM Mode (1 of 2)

The cutoff frequency for the TM<sub>mn</sub> mode was found to be

\[ f_{c,mn} = \frac{1}{2\pi\sqrt{\mu\epsilon}} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2} \]

Note, it is not possible to have \( n = 0 \) or \( m = 0 \) for the TM mode. So...

- The TM<sub>00</sub> mode does not exist.
- The TM<sub>01</sub> mode does not exist.
- The TM<sub>10</sub> mode does not exist.
- The TM<sub>02</sub> mode does not exist.
- The TM<sub>20</sub> mode does not exist.
- The TM<sub>03</sub> mode does not exist.
- The TM<sub>30</sub> mode does not exist.
- etc.

Cutoff for First-Order TM Mode (2 of 2)

What combination of \( m \) and \( n \) minimizes \( f_c \)?

\[ f_{c,mn} = \frac{1}{2\pi\sqrt{\mu\epsilon}} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2} \]

The TM<sub>11</sub> mode will have the lowest cutoff frequency.

\[ m = 1, n = 1 \]

\[ f_{c11} = \frac{1}{2\pi\sqrt{\mu\epsilon}} \sqrt{\left(\frac{1 \cdot \pi}{a}\right)^2 + \left(\frac{1 \cdot \pi}{b}\right)^2} = \frac{1}{2\sqrt{\mu\epsilon}} \sqrt{\left(\frac{1}{a}\right)^2 + \left(\frac{1}{b}\right)^2} \]

CAUTION: It cannot yet be said that the TM<sub>11</sub> is the fundamental mode because the TE modes have not been checked.
Example #2 – TM Mode Analysis (1 of 3)

Given an air-filled rectangular waveguide with \( a = 3 \) cm and \( b = 2 \) cm.

What is the cutoff frequency of the waveguide?

\[
f_{c1} = \frac{1}{2\pi\sqrt{\mu\epsilon}} \left( \frac{\pi}{a} \right)^2 + \left( \frac{\pi}{b} \right)^2
\]

\[
= \frac{1}{2\pi\sqrt{\mu\epsilon\mu_0\epsilon_0}} \left( \frac{\pi}{a} \right)^2 + \left( \frac{\pi}{b} \right)^2
\]

\[
= \frac{1}{2\pi\sqrt{\mu\epsilon}} \sqrt{\left( \frac{\pi}{a} \right)^2 + \left( \frac{\pi}{b} \right)^2}
\]

Recall that \( c_0 = \frac{1}{\sqrt{\mu_0\epsilon_0}} \)

\[
f_{c1} = \frac{29972458 \text{ m/s}}{2\sqrt{(1.0)(1.0)}} \left( \frac{1}{0.03 \text{ m}} \right)^2 + \left( \frac{1}{0.02 \text{ m}} \right)^2 = 9.0 \text{ GHz}
\]

Example #2 – TM Mode Analysis (2 of 3)

Plot the phase constant and effective refractive index for the first-order mode from DC up to 15 GHz.

The phase constant is calculated as:

\[
\beta_m = \sqrt{k^2 - \left( \frac{m\pi}{a} \right)^2 - \left( \frac{n\pi}{b} \right)^2}
\]

The effective refractive index is calculated as:

\[
\beta = k_0 n_{eff} \quad \rightarrow \quad n_{eff} = \frac{\beta}{2\pi f} = \frac{\beta c_0}{2\pi f}
\]
Example #2 – TM Mode Analysis (3 of 3)

Plot the velocity of the modes as a function of frequency.

Are our modes travelling faster than the speed of light?

Summary of TM Analysis

Field Solution

\[ E_x(x, y, z) = -\frac{j\beta_{m,n} \pi}{k_0} B_m \cos \left( \frac{m\pi x}{a} \right) \sin \left( \frac{n\pi y}{b} \right) e^{-j\beta z} \]

\[ E_y(x, y, z) = -\frac{j\beta_{m,n} \pi}{k_0} B_m \sin \left( \frac{m\pi x}{a} \right) \cos \left( \frac{n\pi y}{b} \right) e^{-j\beta z} \]

\[ E_z(x, y, z) = B_m \sin \left( \frac{m\pi x}{a} \right) \sin \left( \frac{n\pi y}{b} \right) e^{-j\beta z} \]

\[ H_x(x, y, z) = \frac{j\omega \mu \pi}{k_0} B_m \sin \left( \frac{m\pi x}{a} \right) \cos \left( \frac{n\pi y}{b} \right) e^{-j\beta z} \]

\[ H_y(x, y, z) = \frac{j\omega \mu \pi}{k_0} B_m \cos \left( \frac{m\pi x}{a} \right) \sin \left( \frac{n\pi y}{b} \right) e^{-j\beta z} \]

\[ H_z(x, y, z) = 0 \]

Phase Constant

\[ \beta_{m,n} = k \sqrt{\left( \frac{m\pi}{a} \right)^2 + \left( \frac{n\pi}{b} \right)^2} \]

Cutoff Frequency

\[ f_{c,m,n} = \frac{1}{2\pi \sqrt{\mu \varepsilon}} \sqrt{\left( \frac{m\pi}{a} \right)^2 + \left( \frac{n\pi}{b} \right)^2} \]

Characteristic Impedance

\[ Z_{TM,m,n} = \eta \frac{\beta_{m,n}}{k} \]

Same equation as for TE

Same equation as for TE

• \( m \neq 0 \) and \( n \neq 0 \), so \( TM_{00} \), \( TM_{01} \), \( TM_{02} \), \( TM_{10} \), \( TM_{20} \), etc. are not supported modes.

• \( TM_{11} \) is the lowest order TM mode
Visualization of Modes

Animation of TE\textsubscript{10}

Notice one bright spot along \(x\) and zero along \(y\). \((m = 1, n = 0)\)
Animation of $\text{TE}_{20}$

Notice two bright spots along $x$ and zero along $y$. ($m = 2, n = 0$)

Animation of $\text{TE}_{01}$

Notice zero bright spots along $x$ and one along $y$. ($m = 0, n = 1$)
Animation of TE\textsubscript{11}

Notice one bright spot along \(x\) and one along \(y\). \((m = 1, n = 1)\)

Animation of TE\textsubscript{21}

Notice two bright spots along \(x\) and one along \(y\). \((m = 2, n = 1)\)
Animation of $TM_{11}$

Conclusions
The Fundamental Mode

The fundamental mode is the mode which has the lowest cutoff frequency. This is either the TE\(_{10}\) or the TM\(_{11}\) mode.

\[
f_{c,\text{TE}} = \frac{1}{2\sqrt{\mu\varepsilon}} \sqrt{\left(\frac{1}{a}\right)^2}
\]

\[
f_{c,\text{TM}} = \frac{1}{2\sqrt{\mu\varepsilon}} \sqrt{\left(\frac{1}{a}\right)^2 + \left(\frac{1}{b}\right)^2}
\]

It can be observed that the TE\(_{10}\) mode will always have the lowest cutoff frequency.

It is concluded that the TE\(_{10}\) mode is the fundamental mode of the waveguide.

This is also called the dominant mode. When multiple modes are excited, usually most of the power ends up in the fundamental mode.

Example #3 – Mode Analysis (1 of 3)

Given an air-filled rectangular waveguide with \(a = 4\) cm and \(b = 2\) cm,

Over what range of frequencies is this waveguide single mode?

An easy way to do this is to calculate a table using a desktop computer.
Given an air-filled rectangular waveguide with $a = 4$ cm and $b = 2$ cm,

Over what range of frequencies is this waveguide single mode?

An easy way to do this is to calculate a table using a desktop computer.

Then sort the table in order of increasing cutoff frequency.

<table>
<thead>
<tr>
<th>Mode</th>
<th>Cutoff</th>
</tr>
</thead>
<tbody>
<tr>
<td>TE10</td>
<td>3.74 GHz</td>
</tr>
<tr>
<td>TE20</td>
<td>7.48 GHz</td>
</tr>
<tr>
<td>TE01</td>
<td>7.48 GHz</td>
</tr>
<tr>
<td>TE11</td>
<td>8.37 GHz</td>
</tr>
<tr>
<td>TM11</td>
<td>8.37 GHz</td>
</tr>
<tr>
<td>TE21</td>
<td>10.58 GHz</td>
</tr>
<tr>
<td>TM21</td>
<td>10.58 GHz</td>
</tr>
<tr>
<td>TE30</td>
<td>11.22 GHz</td>
</tr>
<tr>
<td>TE31</td>
<td>13.49 GHz</td>
</tr>
<tr>
<td>TM31</td>
<td>13.49 GHz</td>
</tr>
<tr>
<td>TE40</td>
<td>14.96 GHz</td>
</tr>
<tr>
<td>TE02</td>
<td>14.96 GHz</td>
</tr>
<tr>
<td>TE12</td>
<td>15.43 GHz</td>
</tr>
<tr>
<td>TM12</td>
<td>15.43 GHz</td>
</tr>
<tr>
<td>TE41</td>
<td>16.73 GHz</td>
</tr>
</tbody>
</table>

It is immediately seen that the TE$_{10}$ mode is the fundamental mode with the lowest cutoff frequency of 3.74 GHz.

The overall range of frequencies for single-mode operation is therefore $3.74 \text{ GHz} < f < 7.48 \text{ GHz}$. 

The second-order mode is taken from the table to be either the TE$_{20}$ or TE$_{01}$ mode because both of these have the same cutoff frequency of 7.48 GHz.
Key Points

- The rectangular waveguide is not a transmission line because it has less than two conductors.
- When filled with a homogeneous dielectric, the rectangular waveguide supports TE and TM modes, but not TEM.
- The cutoff frequencies for $\text{TE}_{mn}$ and $\text{TM}_{mn}$ modes are the same.
- The $\text{TE}_{00}$ mode does not exist.
- For $\text{TM}_{mn}$ modes, $m \neq 0$ and $n \neq 0$.
- The $\text{TE}_{10}$ is the dominant mode because the $\text{TM}_{10}$ mode does not exist.
- Phase velocity of the modes exceeds the vacuum speed of light.

Summary of Rectangular Waveguide

<table>
<thead>
<tr>
<th>Parameter</th>
<th>$\text{TE}_{mn}$ $m = n = 0$ not allowed</th>
<th>$\text{TM}_{mn}$ $m \neq 0$ and $n \neq 0$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$k$</td>
<td>$\alpha \sqrt{\mu}c$</td>
<td>$\alpha \sqrt{\mu}c$</td>
</tr>
<tr>
<td>$k_c$</td>
<td>$\sqrt{(aw/a)^2 - (az/b)^2}$</td>
<td>$\sqrt{(aw/a)^2 - (az/b)^2}$</td>
</tr>
<tr>
<td>$\beta$</td>
<td>$\sqrt{k^2 - k_c^2}$</td>
<td>$\sqrt{k^2 - k_c^2}$</td>
</tr>
<tr>
<td>$\lambda_a$</td>
<td>$2\pi / k_c$</td>
<td>$2\pi / k_c - 2\pi / \alpha$</td>
</tr>
<tr>
<td>$\lambda_p$</td>
<td>$2\pi / \beta_{mn}$</td>
<td>$2\pi / \beta_{mn}$</td>
</tr>
<tr>
<td>$\nu_p$</td>
<td>$\alpha / \beta_{mn}$</td>
<td>$\alpha / \beta_{mn}$</td>
</tr>
<tr>
<td>$\alpha_d$</td>
<td>$k \tan \delta_{2\beta_{mn}}$</td>
<td>$k \tan \delta_{2\beta_{mn}}$</td>
</tr>
<tr>
<td>$E_z$</td>
<td>$\frac{\partial}{\partial x} A_{ze} \cos \left( \alpha x \right)$ $\sin \left( \beta_{mn} y \right)$</td>
<td>$-\frac{\partial}{\partial x} \phi_{zm} \cos \left( \alpha x \right)$ $\sin \left( \beta_{mn} y \right)$</td>
</tr>
<tr>
<td>$E_y$</td>
<td>$\frac{\partial}{\partial y} A_{ye} \cos \left( \beta_{mn} y \right)$ $\sin \left( \alpha x \right)$</td>
<td>$-\frac{\partial}{\partial y} A_{ym} \cos \left( \beta_{mn} y \right)$ $\sin \left( \alpha x \right)$</td>
</tr>
<tr>
<td>$E_x$</td>
<td>0</td>
<td>$\frac{\partial}{\partial z} \phi_{zm} \cos \left( \alpha x \right)$ $\sin \left( \beta_{mn} y \right)$</td>
</tr>
<tr>
<td>$H_x$</td>
<td>$-\frac{\partial}{\partial y} A_{ye} \cos \left( \beta_{mn} y \right)$ $\sin \left( \alpha x \right)$</td>
<td>$\frac{\partial}{\partial y} A_{ym} \cos \left( \beta_{mn} y \right)$ $\sin \left( \alpha x \right)$</td>
</tr>
<tr>
<td>$H_y$</td>
<td>$-\frac{\partial}{\partial z} \phi_{zm} \cos \left( \alpha x \right)$ $\sin \left( \beta_{mn} y \right)$</td>
<td>$\frac{\partial}{\partial z} A_{zm} \cos \left( \alpha x \right)$ $\sin \left( \beta_{mn} y \right)$</td>
</tr>
<tr>
<td>$H_z$</td>
<td>$\frac{\partial}{\partial x} A_{ze} \cos \left( \alpha x \right)$ $\sin \left( \beta_{mn} y \right)$</td>
<td>0</td>
</tr>
<tr>
<td>$Z$</td>
<td>$2k\beta_{mn}$</td>
<td>$\beta_n \eta / k$</td>
</tr>
</tbody>
</table>