



Advanced Computation:
Computational Electromagnetics

Maxwell's Equations in Fourier Space



Outline

- What is Fourier space?
- Complex Fourier series in terms of the reciprocal lattice vectors
- Maxwell's equations in Fourier space
- Visualizing the plane wave expansion

What is Fourier Space?

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Fourier-Space Vs. Frequency-Domain

Fourier transform x , y ,
and z to k_x , k_y , and k_z .

$$\begin{aligned}\nabla \times \vec{E} &= -\mu \frac{\partial \vec{H}}{\partial t} \\ \nabla \times \vec{H} &= \epsilon \frac{\partial \vec{E}}{\partial t}\end{aligned}$$

Fourier transform t to ω .

$$\begin{aligned}j\vec{k} \times \vec{E} &= -\mu \frac{\partial \vec{H}}{\partial t} \\ j\vec{k} \times \vec{H} &= \epsilon \frac{\partial \vec{E}}{\partial t}\end{aligned}$$

Fourier Space

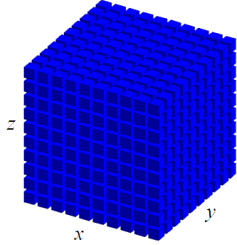
$$\begin{aligned}\nabla \times \vec{E} &= -j\omega\mu\vec{H} \\ \nabla \times \vec{H} &= j\omega\epsilon\vec{E}\end{aligned}$$

Frequency Domain

	Real-Space	Fourier-Space
Time-Domain	FDTD, Discontinuous Galerkin	Pseudo-spectral FDTD
Frequency-Domain	FDFD, FEM, MoM, MoL	RCWA, SAM, PWEM, spectral domain

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What is Fourier Space?

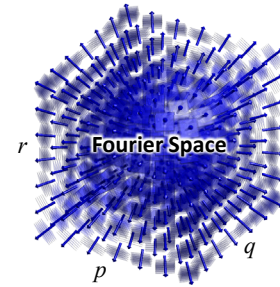


Real Space

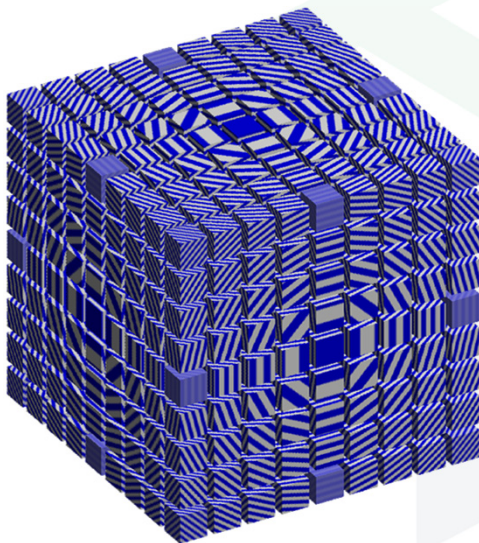
So far, fields and devices were represented on an x - y - z grid where field values and material properties are defined at discrete points.

Fourier Space

In Fourier-space, fields are represented as a sum of plane waves at different angles and different wavelengths called *spatial harmonics*. Devices are also represented as the sum of sinusoidal gratings at different angles and periods.



Visualizing the Spatial Harmonics



$$\vec{k}(p, q, r) = p\vec{T}_1 + q\vec{T}_2 + r\vec{T}_3$$

p , q , and r are the indices of the spatial harmonics.

$$p \equiv \text{integer} \quad -\infty, \dots, -2, -1, 0, 1, 2, \dots, \infty$$

$$q \equiv \text{integer} \quad -\infty, \dots, -2, -1, 0, 1, 2, \dots, \infty$$

$$r \equiv \text{integer} \quad -\infty, \dots, -2, -1, 0, 1, 2, \dots, \infty$$

\vec{T}_1 , \vec{T}_2 , and \vec{T}_3 are the reciprocal lattice vectors.

Each of these plane waves will be assigned its own complex amplitude to convey its magnitude and phase.

Complex Fourier Series in Terms of the Reciprocal Lattice Vectors

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Conventional Complex Fourier Series

Periodic functions can be expanded into a Fourier series.

For 1D periodic functions, this is

$$f(x) = \sum_{p=-\infty}^{\infty} a(p) e^{j \frac{2\pi px}{\Lambda}} \quad a(p) = \frac{1}{\Lambda} \int_{-\Lambda/2}^{\Lambda/2} f(x) e^{-j \frac{2\pi px}{\Lambda}} dx$$

For 2D periodic functions, this is

$$f(x, y) = \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} a(p, q) e^{j \left(\frac{2\pi px}{\Lambda_x} + \frac{2\pi qy}{\Lambda_y} \right)} \quad a(p, q) = \frac{1}{A} \iint_A f(x, y) e^{-j \left(\frac{2\pi px}{\Lambda_x} + \frac{2\pi qy}{\Lambda_y} \right)} dA$$

For 3D periodic functions, this is

$$f(x, y, z) = \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} a(p, q, r) e^{j \left(\frac{2\pi px}{\Lambda_x} + \frac{2\pi qy}{\Lambda_y} + \frac{2\pi rz}{\Lambda_z} \right)} \quad a(p, q, r) = \frac{1}{V} \iiint_V f(x, y, z) e^{-j \left(\frac{2\pi px}{\Lambda_x} + \frac{2\pi qy}{\Lambda_y} + \frac{2\pi rz}{\Lambda_z} \right)} dV$$

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Generalized Complex Fourier Series

Fourier series can be written in terms of the reciprocal lattice vectors.

For 1D periodic functions, this is

$$f(x) = \sum_{p=-\infty}^{\infty} a(p) e^{jpTx} \quad a(p) = \frac{1}{\Lambda} \int_{-\Lambda/2}^{\Lambda/2} f(x) e^{-jpTx} dx \quad T = \frac{2\pi}{\Lambda}$$

For 2D periodic functions, this is

$$f(x, y) = \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} a(p, q) e^{j(p\vec{T}_1 + q\vec{T}_2) \cdot \vec{r}} \quad a(p, q) = \frac{1}{A} \iint_A f(x, y) e^{-j(p\vec{T}_1 + q\vec{T}_2) \cdot \vec{r}} dA$$

For 3D periodic functions, this is

$$f(\vec{r}) = \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} a(p, q, r) e^{j(p\vec{T}_1 + q\vec{T}_2 + r\vec{T}_3) \cdot \vec{r}} \quad a(p, q, r) = \frac{1}{V} \iiint_V f(\vec{r}) e^{-j(p\vec{T}_1 + q\vec{T}_2 + r\vec{T}_3) \cdot \vec{r}} dV$$

For rectangular, tetrahedral, or orthorhombic geometries, the reciprocal lattice vectors are:

$$\vec{T}_1 = \frac{2\pi}{\Lambda_x} \hat{x} \quad \vec{T}_2 = \frac{2\pi}{\Lambda_y} \hat{y} \quad \vec{T}_3 = \frac{2\pi}{\Lambda_z} \hat{z}$$

Maxwell's Equations in Fourier Space

Starting Point

Start with Maxwell's equations in the following form...

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = k_0 \mu_r \tilde{H}_x$$

$$\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} = k_0 \mu_r \tilde{H}_y$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = k_0 \mu_r \tilde{H}_z$$

$$\frac{\partial \tilde{H}_z}{\partial y} - \frac{\partial \tilde{H}_y}{\partial z} = k_0 \epsilon_r E_x$$

$$\frac{\partial \tilde{H}_x}{\partial z} - \frac{\partial \tilde{H}_z}{\partial x} = k_0 \epsilon_r E_y$$

$$\frac{\partial \tilde{H}_y}{\partial x} - \frac{\partial \tilde{H}_x}{\partial y} = k_0 \epsilon_r E_z$$

Recall that the magnetic field was normalized according to

$$\tilde{H} = -j \sqrt{\frac{\mu_0}{\epsilon_0}} \bar{H}$$

Fourier Expansion of the Materials

Assuming the device is infinitely periodic in all directions, the permittivity and permeability functions can be expanded into a generalized Fourier Series.

$$\epsilon_r(\vec{r}) = \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} a(p, q, r) e^{j(p\bar{T}_1 + q\bar{T}_2 + r\bar{T}_3) \cdot \vec{r}}$$

$$a(p, q, r) = \frac{1}{V} \iiint_V \epsilon_r(\vec{r}) e^{-j(p\bar{T}_1 + q\bar{T}_2 + r\bar{T}_3) \cdot \vec{r}} dV$$

$$\mu_r(\vec{r}) = \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} b(p, q, r) e^{j(p\bar{T}_1 + q\bar{T}_2 + r\bar{T}_3) \cdot \vec{r}}$$

$$b(p, q, r) = \frac{1}{V} \iiint_V \mu_r(\vec{r}) e^{-j(p\bar{T}_1 + q\bar{T}_2 + r\bar{T}_3) \cdot \vec{r}} dV$$

Fourier Expansion of the Fields

The expansions are slightly different for fields because a wave could be travelling in any direction $\vec{\beta}$. The field must obey Bloch's theorem.

$$\vec{E}(\vec{r}) = e^{-j\vec{\beta}\cdot\vec{r}} \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} \vec{S}(p, q, r) e^{j(p\vec{T}_1 + q\vec{T}_2 + r\vec{T}_3)\cdot\vec{r}}$$

Think of $\vec{\beta}$ as \vec{k}_{inc} , but be careful with that analogy.

$$= \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} \vec{S}(p, q, r) e^{-j(\vec{\beta} - p\vec{T}_1 - q\vec{T}_2 - r\vec{T}_3)\cdot\vec{r}}$$

Let this be $\vec{k}(p, q, r)$

$e^{-j\vec{\beta}\cdot\vec{r}}$ was brought inside summation and combined with second exponential.

$$= \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} \vec{S}(p, q, r) e^{-j\vec{k}(p, q, r)\cdot\vec{r}}$$

This is clearly a set of plane waves with amplitudes $\vec{S}(p, q, r)$.

$$= \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} \vec{S}(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]}$$

$$\vec{k}(p, q, r) = \vec{\beta} - p\vec{T}_1 - q\vec{T}_2 - r\vec{T}_3$$

$$k_x(p, q, r) = \beta_x - pT_{1,x} - qT_{2,x} - rT_{3,x}$$

$$k_y(p, q, r) = \beta_y - pT_{1,y} - qT_{2,y} - rT_{3,y}$$

$$k_z(p, q, r) = \beta_z - pT_{1,z} - qT_{2,z} - rT_{3,z}$$

Substitute Expansions into Maxwell's Equations

$$\vec{H}_y(\vec{r}) = \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_y(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]}$$

$$\vec{H}_z(\vec{r}) = \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_z(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]}$$

$$\epsilon_r(\vec{r}) = \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} a(p, q, r) e^{j(p\vec{T}_1 + q\vec{T}_2 + r\vec{T}_3)\cdot\vec{r}}$$

$$E_x(\vec{r}) = \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} S_x(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]}$$

$$\frac{\partial \vec{H}_z}{\partial y} - \frac{\partial \vec{H}_y}{\partial z} = k_0 \epsilon_r E_x$$

$$\frac{\partial}{\partial y} \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_z(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \right] - \frac{\partial}{\partial z} \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_y(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \right]$$

$$= k_0 \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} a(p, q, r) e^{j(p\vec{T}_1 + q\vec{T}_2 + r\vec{T}_3)\cdot\vec{r}} \right] \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} S_x(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \right]$$

Algebra for the Left Side Terms

First ugly term...

$$\begin{aligned} \frac{\partial}{\partial y} \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_z(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \right] &= \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_z(p, q, r) \frac{\partial}{\partial y} e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \\ &= \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_z(p, q, r) [-jk_y(p, q, r)] e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \\ &= \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} -jk_y(p, q, r) U_z(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \end{aligned}$$

Second ugly term...

$$\begin{aligned} \frac{\partial}{\partial z} \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_y(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \right] &= \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_y(p, q, r) \frac{\partial}{\partial z} e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \\ &= \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_y(p, q, r) [-jk_z(p, q, r)] e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \\ &= \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} -jk_z(p, q, r) U_y(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \end{aligned}$$

Algebra for the Right-Side Term

Third ugly term...

This term as the product of two triple summations.

$$\left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} a(p, q, r) e^{j(p\vec{T}_1 + q\vec{T}_2 + r\vec{T}_3) \cdot \vec{r}} \right] \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} S_x(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \right]$$

This is called a Cauchy product and is handled as follows.

$$\left(\sum_{n=0}^{\infty} a_n \right) \cdot \left(\sum_{n=0}^{\infty} b_n \right) = \sum_{n=0}^{\infty} c_n \quad c_n = \sum_{m=0}^n a_m b_{n-m}$$

Applying this rule to the triple summations, gives

$$\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} \left\{ e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \sum_{p'=-\infty}^{\infty} \sum_{q'=-\infty}^{\infty} \sum_{r'=-\infty}^{\infty} a(p-p', q-q', r-r') S_x(p', q', r') \right\}$$

Combine the Terms Inside Summation

$$\frac{\partial}{\partial y} \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_z(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \right] - \frac{\partial}{\partial z} \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} U_y(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \right]$$

$$= k_0 \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} a(p, q, r) e^{j(p\bar{t}_1 + q\bar{t}_2 + r\bar{t}_3) \cdot \bar{r}} \right] \left[\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} S_x(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \right]$$



$$\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} -jk_y(p, q, r) U_z(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} - \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} -jk_z(p, q, r) U_y(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]}$$

$$= k_0 \sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} \left\{ e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \sum_{p'=-\infty}^{\infty} \sum_{q'=-\infty}^{\infty} \sum_{r'=-\infty}^{\infty} a(p-p', q-q', r-r') S_x(p', q', r') \right\}$$

The equation can now be brought inside a single triple summation.

$$\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} \left\{ \begin{aligned} & -jk_y(p, q, r) U_z(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} + jk_z(p, q, r) U_y(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \\ & = k_0 e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \sum_{p'=-\infty}^{\infty} \sum_{q'=-\infty}^{\infty} \sum_{r'=-\infty}^{\infty} a(p-p', q-q', r-r') S_x(p', q', r') \end{aligned} \right\}$$

Final Equation for $(p, q, r)^{\text{th}}$ Harmonic

$$\sum_{p=-\infty}^{\infty} \sum_{q=-\infty}^{\infty} \sum_{r=-\infty}^{\infty} \left\{ \begin{aligned} & -jk_y(p, q, r) U_z(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} + jk_z(p, q, r) U_y(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \\ & = k_0 e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \sum_{p'=-\infty}^{\infty} \sum_{q'=-\infty}^{\infty} \sum_{r'=-\infty}^{\infty} a(p-p', q-q', r-r') S_x(p', q', r') \end{aligned} \right\}$$

The equation inside the braces must be satisfied for each combination of (p, q, r) .

$$-jk_y(p, q, r) U_z(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} + jk_z(p, q, r) U_y(p, q, r) e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]}$$

$$= k_0 e^{-j[k_x(p, q, r)x + k_y(p, q, r)y + k_z(p, q, r)z]} \sum_{p'=-\infty}^{\infty} \sum_{q'=-\infty}^{\infty} \sum_{r'=-\infty}^{\infty} a(p-p', q-q', r-r') S_x(p', q', r')$$

Last, divide both sides by the common exponential term and move the j to the right-hand side.

$$k_y(p, q, r) U_z(p, q, r) - k_z(p, q, r) U_y(p, q, r) = jk_0 \sum_{p'=-\infty}^{\infty} \sum_{q'=-\infty}^{\infty} \sum_{r'=-\infty}^{\infty} a(p-p', q-q', r-r') S_x(p', q', r')$$

Alternate Derivation

Start with

$$\frac{\partial \tilde{H}_z}{\partial y} - \frac{\partial \tilde{H}_y}{\partial z} = k_0 \epsilon_r E_x$$

Point-by-point multiplication in real-space...

Fourier-transform this equation in x , y , and z resulting in

$$k_y(p, q, r)U_z(p, q, r) - k_z(p, q, r)U_y(p, q, r) = jk_0 a * S_x$$

...becomes a convolution in Fourier-space.

$a = \text{FT}\{\epsilon_r\}$
 $S_x = \text{FT}\{E_x\}$

It can now be seen that the strange triple summation remaining in the Fourier-space equation is actually a 3D convolution in Fourier space!

$$a * S_x \rightarrow \sum_{p'=-\infty}^{\infty} \sum_{q'=-\infty}^{\infty} \sum_{r'=-\infty}^{\infty} a(p-p', q-q', r-r') S_x(p', q', r')$$

Summary of Maxwell's Equations in Fourier Space

Real-Space

$$\frac{\partial \tilde{H}_z}{\partial y} - \frac{\partial \tilde{H}_y}{\partial z} = k_0 \epsilon_r E_x$$

$$\frac{\partial \tilde{H}_x}{\partial z} - \frac{\partial \tilde{H}_z}{\partial x} = k_0 \epsilon_r E_y$$

$$\frac{\partial \tilde{H}_y}{\partial x} - \frac{\partial \tilde{H}_x}{\partial y} = k_0 \epsilon_r E_z$$

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = k_0 \mu_r \tilde{H}_x$$

$$\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} = k_0 \mu_r \tilde{H}_y$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = k_0 \mu_r \tilde{H}_z$$

Fourier-Space

$$\begin{aligned} k_y(p, q, r)U_z(p, q, r) - k_z(p, q, r)U_y(p, q, r) &= jk_0 a(p, q, r) * S_x(p, q, r) \\ k_z(p, q, r)U_x(p, q, r) - k_x(p, q, r)U_z(p, q, r) &= jk_0 a(p, q, r) * S_y(p, q, r) \\ k_x(p, q, r)U_y(p, q, r) - k_y(p, q, r)U_x(p, q, r) &= jk_0 a(p, q, r) * S_z(p, q, r) \end{aligned}$$

$$\vec{k}(p, q, r) = k_x(p, q, r)\hat{a}_x + k_y(p, q, r)\hat{a}_y + k_z(p, q, r)\hat{a}_z = \vec{\beta} - p\vec{T}_1 - q\vec{T}_2 - r\vec{T}_3$$

$$p = -\infty, \dots, -2, -1, 0, 1, 2, \dots, \infty$$

$$q = -\infty, \dots, -2, -1, 0, 1, 2, \dots, \infty$$

$$r = -\infty, \dots, -2, -1, 0, 1, 2, \dots, \infty$$

$$\begin{aligned} k_y(p, q, r)S_z(p, q, r) - k_z(p, q, r)S_y(p, q, r) &= jk_0 b(p, q, r) * U_x(p, q, r) \\ k_z(p, q, r)S_x(p, q, r) - k_x(p, q, r)S_z(p, q, r) &= jk_0 b(p, q, r) * U_y(p, q, r) \\ k_x(p, q, r)S_y(p, q, r) - k_y(p, q, r)S_x(p, q, r) &= jk_0 b(p, q, r) * U_z(p, q, r) \end{aligned}$$

Vector Form of Maxwell's Equations in Fourier Space

$$k_y(p, q, r)U_z(p, q, r) - k_z(p, q, r)U_y(p, q, r) = jk_0 a(p, q, r) * S_x(p, q, r)$$

$$k_z(p, q, r)U_x(p, q, r) - k_x(p, q, r)U_z(p, q, r) = jk_0 a(p, q, r) * S_y(p, q, r)$$

$$k_x(p, q, r)U_y(p, q, r) - k_y(p, q, r)U_x(p, q, r) = jk_0 a(p, q, r) * S_z(p, q, r)$$

$$\vec{k}(p, q, r) \times \vec{U}(p, q, r) = jk_0 \epsilon_r(p, q, r) * \vec{S}(p, q, r)$$

$$k_y(p, q, r)S_z(p, q, r) - k_z(p, q, r)S_y(p, q, r) = jk_0 b(p, q, r) * U_x(p, q, r)$$

$$k_z(p, q, r)S_x(p, q, r) - k_x(p, q, r)S_z(p, q, r) = jk_0 b(p, q, r) * U_y(p, q, r)$$

$$k_x(p, q, r)S_y(p, q, r) - k_y(p, q, r)S_x(p, q, r) = jk_0 b(p, q, r) * U_z(p, q, r)$$

$$\vec{k}(p, q, r) \times \vec{S}(p, q, r) = jk_0 \mu_r(p, q, r) * \vec{U}(p, q, r)$$

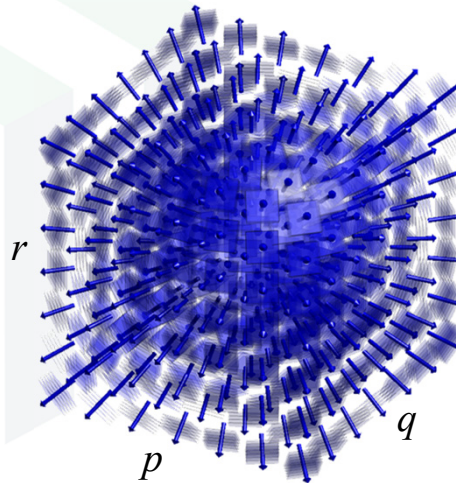
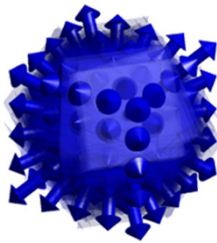
Visualizing the Plane Wave Expansion

Visualizing Maxwell's Equations in Fourier Space

In real-space, the field values are known at discrete points.

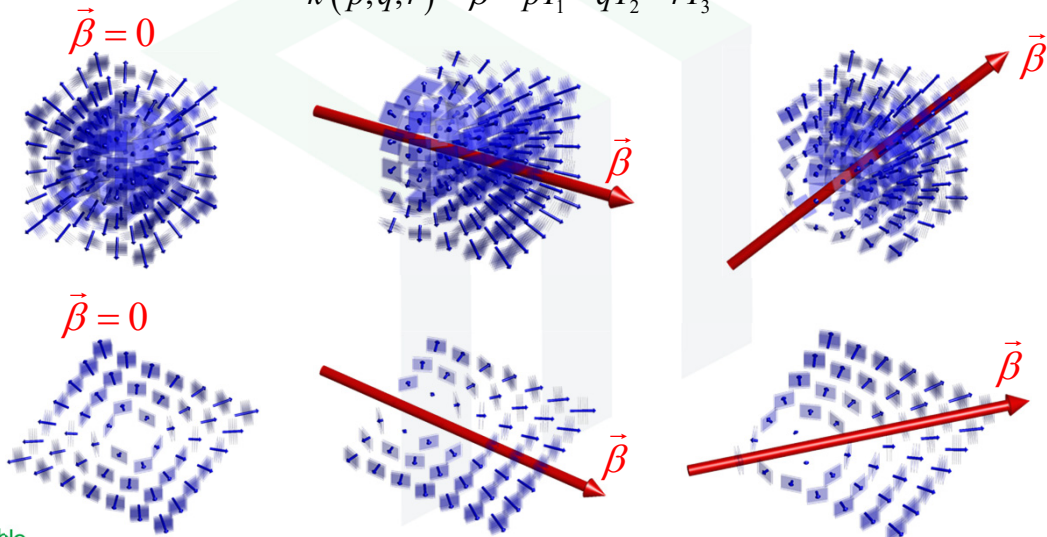
In Fourier-space, amplitudes are known of discrete plane waves.

A less clear, but more accurate picture is when all the plane waves overlap.



Visualizing Expansions for Cubic Symmetry

$$\vec{k}(p, q, r) = \vec{\beta} - p\vec{T}_1 - q\vec{T}_2 - r\vec{T}_3$$



Visualizing Expansions with Different Symmetries

$$\vec{k}(p, q, r) = \vec{\beta} - p\vec{T}_1 - q\vec{T}_2 - r\vec{T}_3$$

