



Advanced Computation:
Computational Electromagnetics

Finite-Difference Approximation of Maxwell's Equations on a Yee Grid



Outline

- The Starting Point
- Finite-Difference Approximation of $\nabla \times \vec{E} = \mu_r \vec{H}$
- Finite-Difference Approximation of $\nabla \times \vec{H} = \epsilon_r \vec{E}$
- Summary of Finite-Difference Equations on a Yee Grid

The Starting Point

Slide 3

Normalize the Grid Coordinates x , y and z

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = k_0 \mu_{xx} \tilde{H}_x$$

$$\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} = k_0 \mu_{yy} \tilde{H}_y$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = k_0 \mu_{zz} \tilde{H}_z$$

$$\frac{\partial \tilde{H}_z}{\partial y} - \frac{\partial \tilde{H}_y}{\partial z} = k_0 \epsilon_{xx} E_x$$

$$\frac{\partial \tilde{H}_x}{\partial z} - \frac{\partial \tilde{H}_z}{\partial x} = k_0 \epsilon_{yy} E_y$$

$$\frac{\partial \tilde{H}_y}{\partial x} - \frac{\partial \tilde{H}_x}{\partial y} = k_0 \epsilon_{zz} E_z$$

Normalize
the grid

$$x' = k_0 x$$

$$y' = k_0 y$$

$$z' = k_0 z$$

This "absorbs" the k_0 term into the spatial derivatives and simplifies Maxwell's equations.

$$\frac{\partial E_z}{\partial y'} - \frac{\partial E_y}{\partial z'} = \mu_{xx} \tilde{H}_x$$

$$\frac{\partial E_x}{\partial z'} - \frac{\partial E_z}{\partial x'} = \mu_{yy} \tilde{H}_y$$

$$\frac{\partial E_y}{\partial x'} - \frac{\partial E_x}{\partial y'} = \mu_{zz} \tilde{H}_z$$

$$\frac{\partial \tilde{H}_z}{\partial y'} - \frac{\partial \tilde{H}_y}{\partial z'} = \epsilon_{xx} E_x$$

$$\frac{\partial \tilde{H}_x}{\partial z'} - \frac{\partial \tilde{H}_z}{\partial x'} = \epsilon_{yy} E_y$$

$$\frac{\partial \tilde{H}_y}{\partial x'} - \frac{\partial \tilde{H}_x}{\partial y'} = \epsilon_{zz} E_z$$

The Starting Point

$$\frac{\partial E_z}{\partial y'} - \frac{\partial E_y}{\partial z'} = \mu_{xx} \tilde{H}_x$$

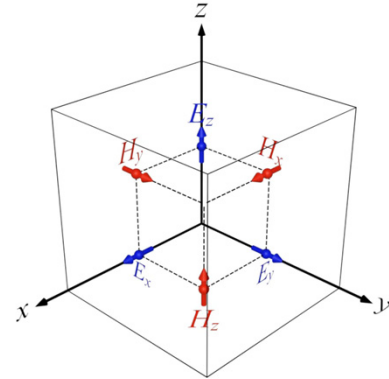
$$\frac{\partial E_x}{\partial z'} - \frac{\partial E_z}{\partial x'} = \mu_{yy} \tilde{H}_y$$

$$\frac{\partial E_y}{\partial x'} - \frac{\partial E_x}{\partial y'} = \mu_{zz} \tilde{H}_z$$

$$\frac{\partial \tilde{H}_z}{\partial y'} - \frac{\partial \tilde{H}_y}{\partial z'} = \varepsilon_{xx} E_x$$

$$\frac{\partial \tilde{H}_x}{\partial z'} - \frac{\partial \tilde{H}_z}{\partial x'} = \varepsilon_{yy} E_y$$

$$\frac{\partial \tilde{H}_y}{\partial x'} - \frac{\partial \tilde{H}_x}{\partial y'} = \varepsilon_{zz} E_z$$



Here, only diagonally anisotropic material tensors were retained. This will be needed to incorporate a uniaxial perfectly matched layer (UPML) absorbing boundary condition.

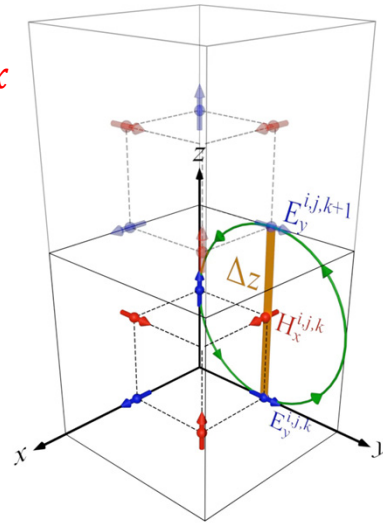
These equations are valid independent of the chosen sign convention.

Finite-Difference Approximation of

$$\nabla \times \vec{E} = \mu_r \vec{H}$$

Finite-Difference Equation for H_x

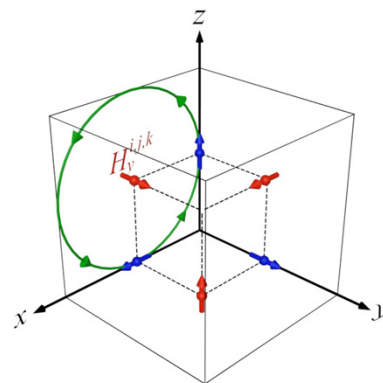
$$\frac{\partial E_z}{\partial y'} - \frac{\partial E_y}{\partial z'} = \mu_{xx} \tilde{H}_x$$



$$\frac{E_z^{i,j+1,k} - E_z^{i,j,k}}{\Delta y'} - \frac{E_y^{i,j,k+1} - E_y^{i,j,k}}{\Delta z'} = \mu_{xx}^{i,j,k} \tilde{H}_x^{i,j,k}$$

Finite-Difference Equation for H_y

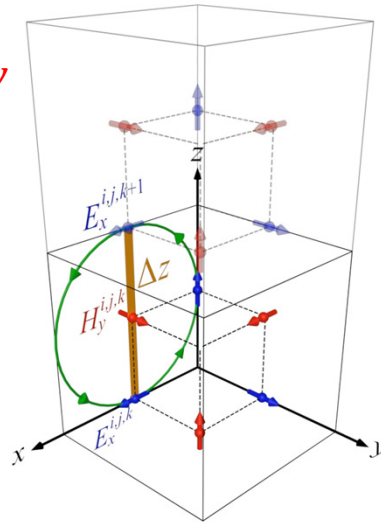
$$\frac{\partial E_x}{\partial z'} - \frac{\partial E_z}{\partial x'} = \mu_{yy} \tilde{H}_y$$



$$= \mu_{yy}^{i,j,k} \tilde{H}_y^{i,j,k}$$

Finite-Difference Equation for H_y

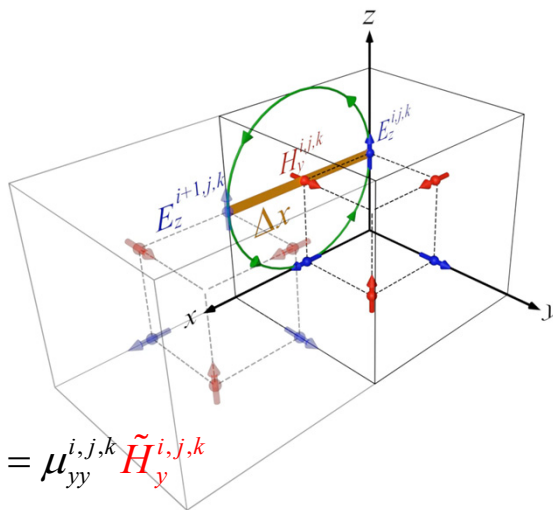
$$\boxed{\frac{\partial E_x}{\partial z'}} - \frac{\partial E_z}{\partial x'} = \mu_{yy} \tilde{H}_y$$



$$\frac{E_x^{i,j,k+1} - E_x^{i,j,k}}{\Delta z'} = \mu_{yy}^{i,j,k} \tilde{H}_y^{i,j,k}$$

Finite-Difference Equation for H_y

$$\frac{\partial E_x}{\partial z'} - \boxed{\frac{\partial E_z}{\partial x'}} = \mu_{yy} \tilde{H}_y$$

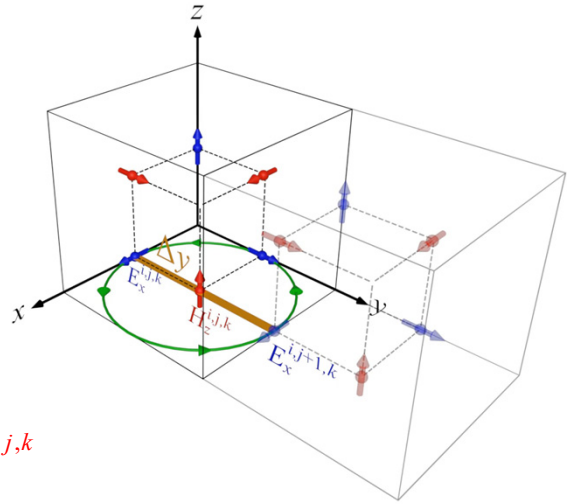


$$\frac{E_x^{i,j,k+1} - E_x^{i,j,k}}{\Delta z'} - \frac{E_z^{i+1,j,k} - E_z^{i,j,k}}{\Delta x'} = \mu_{yy}^{i,j,k} \tilde{H}_y^{i,j,k}$$

Finite-Difference Equation for H_z

$$\frac{\partial E_y}{\partial x'} - \frac{\partial E_x}{\partial y'} = \mu_{zz} \tilde{H}_z$$

$$\frac{E_y^{i+1,j,k} - E_y^{i,j,k}}{\Delta x'} - \frac{E_x^{i,j+1,k} - E_x^{i,j,k}}{\Delta y'} = \mu_{zz}^{i,j,k} \tilde{H}_z^{i,j,k}$$

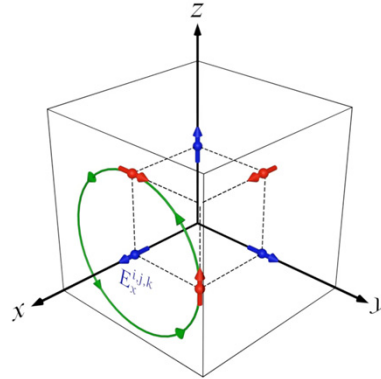


Finite-Difference Approximation of

$$\nabla \times \vec{H} = \epsilon_r \vec{E}$$

Finite-Difference Equation for E_x

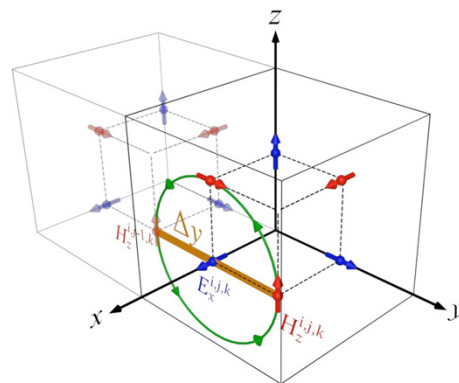
$$\frac{\partial \tilde{H}_z}{\partial y'} - \frac{\partial \tilde{H}_y}{\partial z'} = \epsilon_{xx} E_x$$



$$= \epsilon_{xx}^{i,j,k} E_x^{i,j,k}$$

Finite-Difference Equation for E_x

$$\frac{\partial \tilde{H}_z}{\partial y'} - \frac{\partial \tilde{H}_y}{\partial z'} = \epsilon_{xx} E_x$$



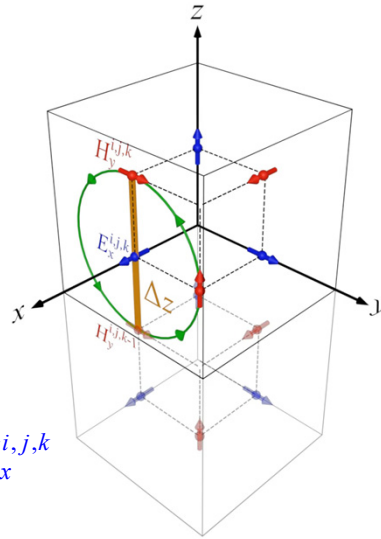
$$\frac{\tilde{H}_z^{i,j,k} - \tilde{H}_z^{i,j-1,k}}{\Delta y'}$$

$$= \epsilon_{xx}^{i,j,k} E_x^{i,j,k}$$

Finite-Difference Equation for E_x

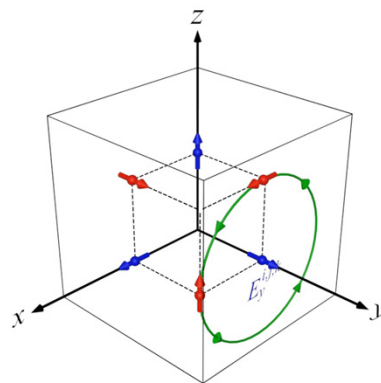
$$\frac{\partial \tilde{H}_z}{\partial y'} - \frac{\partial \tilde{H}_y}{\partial z'} = \epsilon_{xx} E_x$$

$$\frac{\tilde{H}_z^{i,j,k} - \tilde{H}_z^{i,j-1,k}}{\Delta y'} - \frac{\tilde{H}_y^{i,j,k} - \tilde{H}_y^{i,j,k-1}}{\Delta z'} = \epsilon_{xx}^{i,j,k} E_x^{i,j,k}$$

Finite-Difference Equation for E_y

$$\frac{\partial \tilde{H}_x}{\partial z'} - \frac{\partial \tilde{H}_z}{\partial x'} = \epsilon_{yy} E_y$$

$$= \epsilon_{yy}^{i,j,k} E_y^{i,j,k}$$

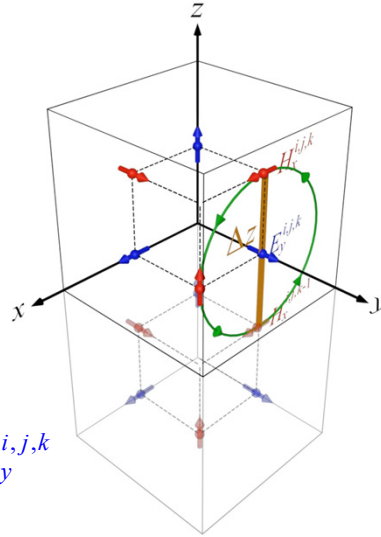


Finite-Difference Equation for E_y

$$\boxed{\frac{\partial \tilde{H}_x}{\partial z'}} - \frac{\partial \tilde{H}_z}{\partial x'} = \epsilon_{yy} E_y$$

$$\frac{\tilde{H}_x^{i,j,k} - \tilde{H}_x^{i,j,k-1}}{\Delta z'}$$

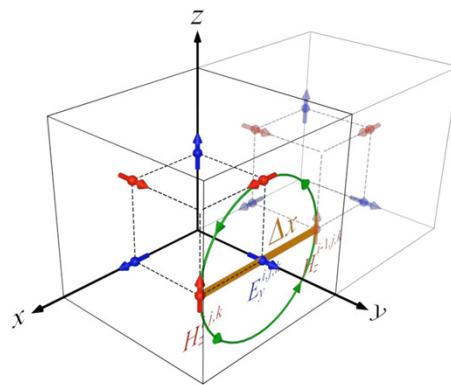
$$= \epsilon_{yy}^{i,j,k} E_y^{i,j,k}$$



Finite-Difference Equation for E_y

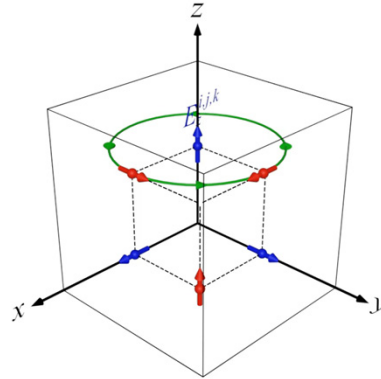
$$\frac{\partial \tilde{H}_x}{\partial z'} - \boxed{\frac{\partial \tilde{H}_z}{\partial x'}} = \epsilon_{yy} E_y$$

$$\frac{\tilde{H}_x^{i,j,k} - \tilde{H}_x^{i,j,k-1}}{\Delta z'} - \frac{\tilde{H}_z^{i,j,k} - \tilde{H}_z^{i-1,j,k}}{\Delta x'} = \epsilon_{yy}^{i,j,k} E_y^{i,j,k}$$



Finite-Difference Equation for E_z

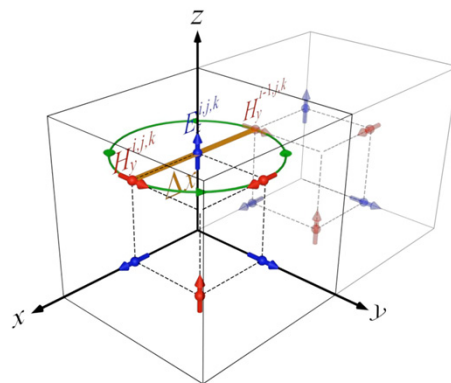
$$\frac{\partial \tilde{H}_y}{\partial x'} - \frac{\partial \tilde{H}_x}{\partial y'} = \epsilon_{zz} E_z$$



$$= \epsilon_{zz}^{i,j,k} E_z^{i,j,k}$$

Finite-Difference Equation for E_z

$$\frac{\partial \tilde{H}_y}{\partial x'} - \frac{\partial \tilde{H}_x}{\partial y'} = \epsilon_{zz} E_z$$



$$\frac{\tilde{H}_y^{i,j,k} - \tilde{H}_y^{i-1,j,k}}{\Delta x'}$$

$$= \epsilon_{zz}^{i,j,k} E_z^{i,j,k}$$

Summary of Finite-Difference Approximations of Maxwell's Equations

$$\begin{array}{l}
 \frac{\partial E_z}{\partial y'} - \frac{\partial E_y}{\partial z'} = \mu_{xx} \tilde{H}_x \\
 \frac{\partial E_x}{\partial z'} - \frac{\partial E_z}{\partial x'} = \mu_{yy} \tilde{H}_y \\
 \frac{\partial E_y}{\partial x'} - \frac{\partial E_x}{\partial y'} = \mu_{zz} \tilde{H}_z \\
 \\
 \frac{\partial \tilde{H}_z}{\partial y'} - \frac{\partial \tilde{H}_y}{\partial z'} = \varepsilon_{xx} E_x \\
 \frac{\partial \tilde{H}_x}{\partial z'} - \frac{\partial \tilde{H}_z}{\partial x'} = \varepsilon_{yy} E_y \\
 \frac{\partial \tilde{H}_y}{\partial x'} - \frac{\partial \tilde{H}_x}{\partial y'} = \varepsilon_{zz} E_z
 \end{array}
 \quad \longrightarrow \quad
 \begin{array}{l}
 \frac{E_z^{i,j+1,k} - E_z^{i,j,k}}{\Delta y'} - \frac{E_y^{i,j,k+1} - E_y^{i,j,k}}{\Delta z'} = \mu_{xx}^{i,j,k} \tilde{H}_x^{i,j,k} \\
 \frac{E_x^{i,j,k+1} - E_x^{i,j,k}}{\Delta z'} - \frac{E_z^{i+1,j,k} - E_z^{i,j,k}}{\Delta x'} = \mu_{yy}^{i,j,k} \tilde{H}_y^{i,j,k} \\
 \frac{E_y^{i+1,j,k} - E_y^{i,j,k}}{\Delta x'} - \frac{E_x^{i,j+1,k} - E_x^{i,j,k}}{\Delta y'} = \mu_{zz}^{i,j,k} \tilde{H}_z^{i,j,k} \\
 \\
 \frac{\tilde{H}_z^{i,j,k} - \tilde{H}_z^{i,j-1,k}}{\Delta y'} - \frac{\tilde{H}_y^{i,j,k} - \tilde{H}_y^{i,j,k-1}}{\Delta z'} = \varepsilon_{xx}^{i,j,k} E_x^{i,j,k} \\
 \frac{\tilde{H}_x^{i,j,k} - \tilde{H}_x^{i,j,k-1}}{\Delta z'} - \frac{\tilde{H}_z^{i,j,k} - \tilde{H}_z^{i-1,j,k}}{\Delta x'} = \varepsilon_{yy}^{i,j,k} E_y^{i,j,k} \\
 \frac{\tilde{H}_y^{i,j,k} - \tilde{H}_y^{i-1,j,k}}{\Delta x'} - \frac{\tilde{H}_x^{i,j,k} - \tilde{H}_x^{i,j-1,k}}{\Delta y'} = \varepsilon_{zz}^{i,j,k} E_z^{i,j,k}
 \end{array}$$